



# Exciting Baryons

with MAMI and MAID

Lothar Tiator (Mainz)



Nucleon Resonances: From Photoproduction to High Photon Virtualities  
Trento, October, 12-16, 2015



# The Roper Resonance



first baryon resonance discovered in PWA, 1964

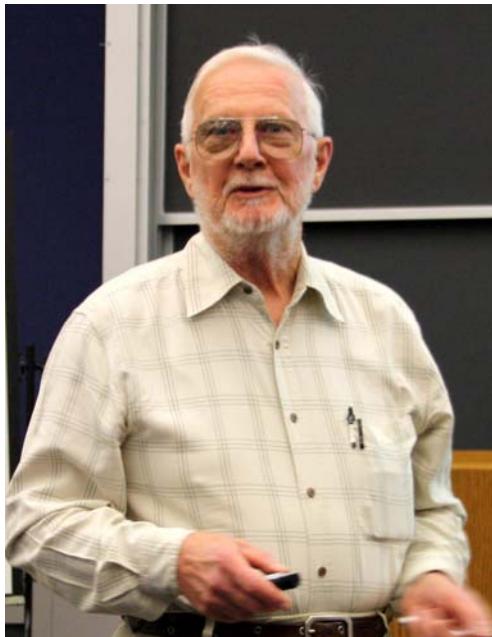
Phys. Rev. Lett. 12 (1964) 340:

EVIDENCE FOR A  $P_{11}$  PION-NUCLEON RESONANCE AT 556 MeV<sup>†</sup>

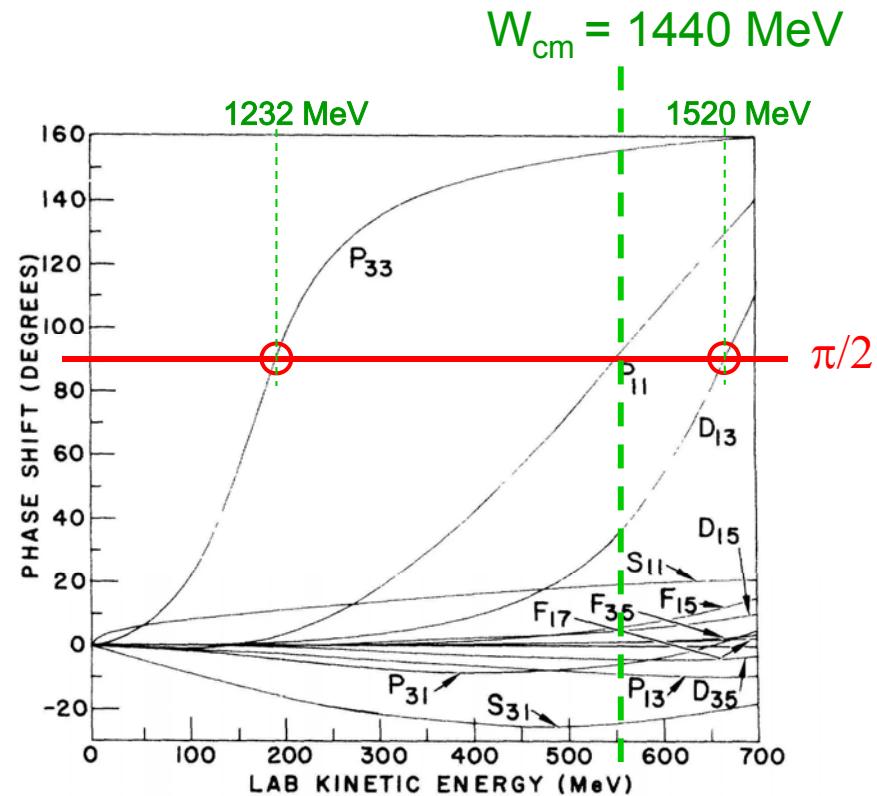
L. David Roper

Lawrence Radiation Laboratory, University of California, Livermore, California

(Received 17 February 1964)



Dan Roper at GWU in 2011



pion-nucleon phase shifts  
from 0 to 700 MeV lab kinetic energy



# outline of my talk

- new MAMI data on photoproduction of  $\pi^0$ ,  $\eta$ ,  $\eta'$  with very high precision cross sections and new polarization observables
- update of EtaMAID for  $\gamma, \eta$  and  $\gamma, \eta'$  with new  $N^*$  excitations and new data from Mainz, Bonn, GRAAL, Jlab
- current status of MAID for  $\gamma, \pi$
- transition form factors from low to medium-high  $Q^2$
- outlook



# collaboration

**Mainz:** Michael Ostrick, Viktor Kashevarov, Kirill Nikonov  
and MAMI-A2 collaboration

**Zagreb:** Alfred Svarc, Sasa Ceci

**Tuzla:** Jugoslav Stahov, Hedim Osmanovic,  
Mirza Hadzimehmedovic

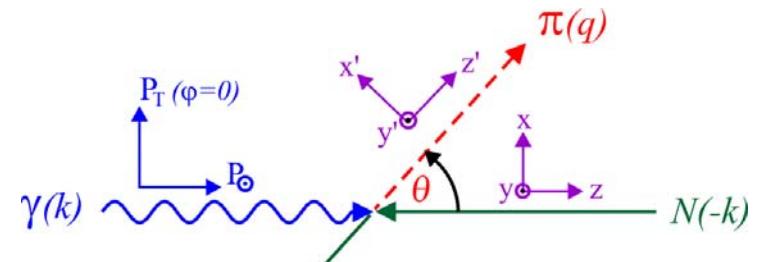
**Dubna:** Sabit Kamalov

# 16 spin observables in photoproduction of $\pi$ , $\eta$ , $\eta'$ , K



linear and circular polarized beams

longitudinal and transverse polarized targets



Photon		Beam - Target		
		-	-	-
		-	x	y
unpolarized	$\sigma$	0	$T$	0
linear polariz.	$\Sigma$	$H$	$P$	$G$
circular polariz.	0	$F$	0	$E$



# recent MAMI data on single $\pi^0$ , $\eta$ , $\eta'$ photoproduction

## on proton:

$\gamma, \pi^0$ :	$d\sigma/d\Omega$ , $\sigma_{total}$	$thr < W < 1.9 \text{ GeV}$	<i>PRL 111 (2013), PRC 92 (2015)</i>
	$T, F$	$thr < W < 1.9 \text{ GeV}$	<i>PL B750 and to be published in 2015</i>
	$G$	$thr < W < 1.55 \text{ GeV}$	<i>not finally analyzed</i>
	$E, \Delta\sigma_{total}$	$thr < W < 1.9 \text{ GeV}$	<i>not finally analyzed</i>
$\gamma, \eta$ :	$d\sigma/d\Omega$ , $\sigma_{total}$	$thr < W < 1.9 \text{ GeV}$	<i>PRC 82 (2010)</i>
		$thr < W < 1.96 \text{ GeV}$	<i>to be published in 2015</i>
	$T, F$	$thr < W < 1.85 \text{ GeV}$	<i>PRL 113 (2014)</i>
	$E, \Delta\sigma_{total}$	$thr < W < 1.9 \text{ GeV}$	<i>not finally analyzed</i>
$\gamma, \eta'$ :	$d\sigma/d\Omega$ , $\sigma_{total}$	$thr < W < 1.96 \text{ GeV}$	<i>to be published in 2015</i>

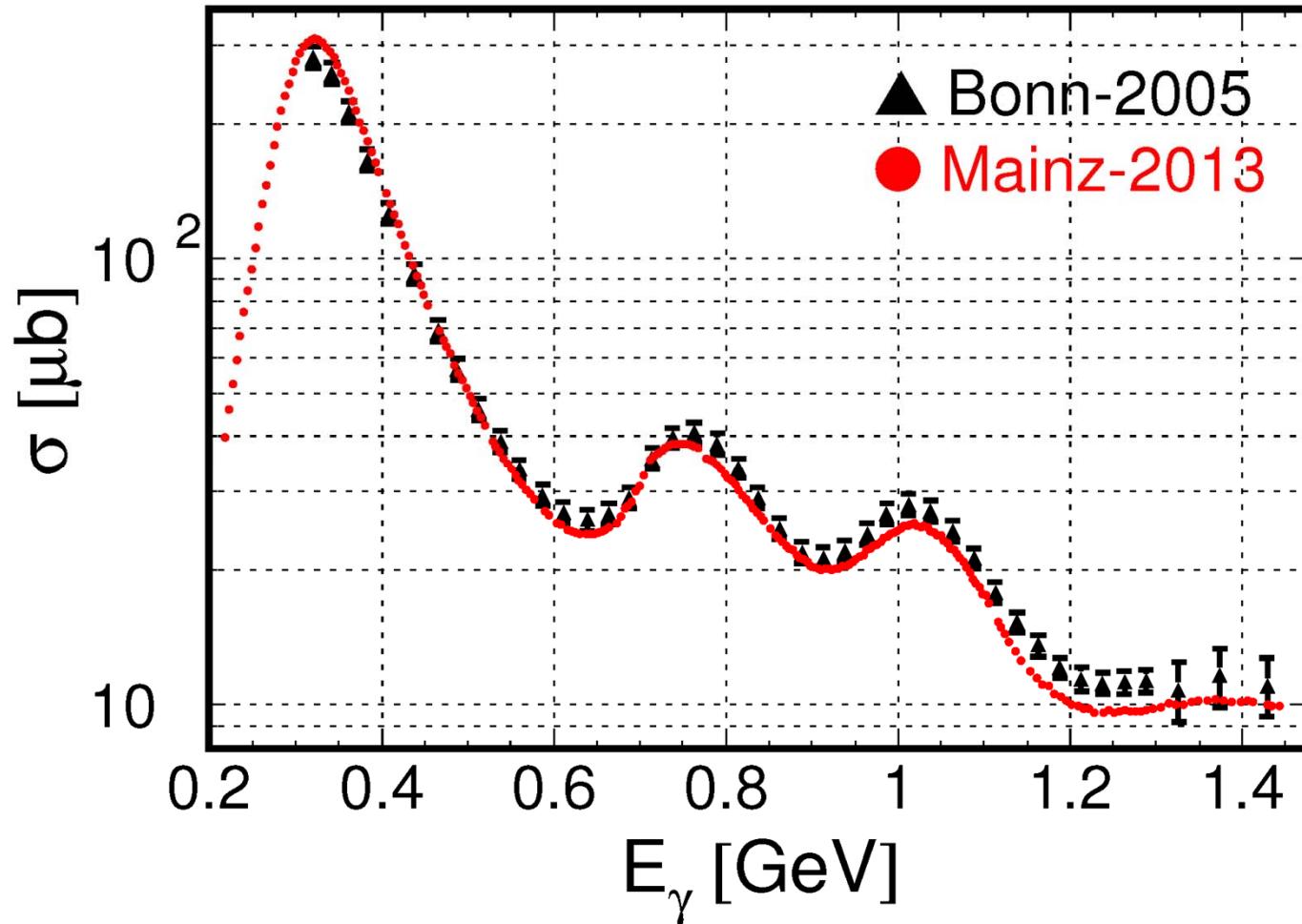
## on neutron:

$\gamma, \pi^0$ :	$d\sigma/d\Omega$ , $\sigma_{total}$	$thr < W < 1.9 \text{ GeV}$	<i>PRL 112 (2014)</i>
$\gamma, \eta$ :	$d\sigma/d\Omega$ , $\sigma_{total}$	$thr < W < 1.9 \text{ GeV}$	<i>PRC 90 (2014)</i>
<i>also</i>	$T, F, E, G$	<i>up to 1.9 GeV</i>	<i>in analysis, to be published in 2016</i>

# total cross section of MAMI $\gamma, \pi^0$ data



P. Adlarson et al (MAMI A2 Collab.) Phys. Rev. C92 (2015) 2, 024617

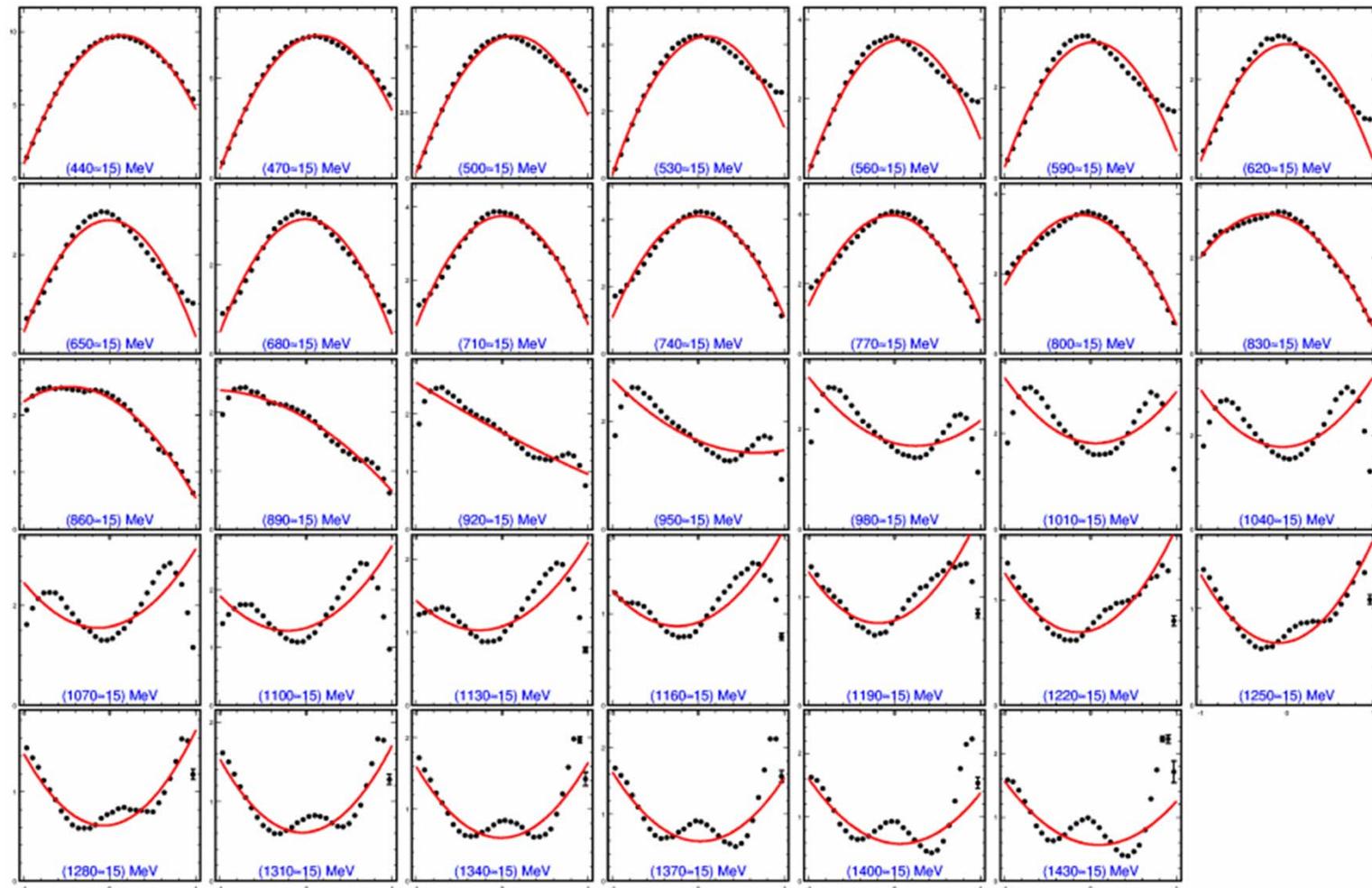


the red points are not a calculation, it is data!

# Legendre expansion of the diff. cross section



$$\frac{d\sigma}{d\Omega} = \sum_{k=0}^{2\ell_{max}} A_k^\sigma(W) P_k(\cos \theta) \quad \ell_{max} = 1$$

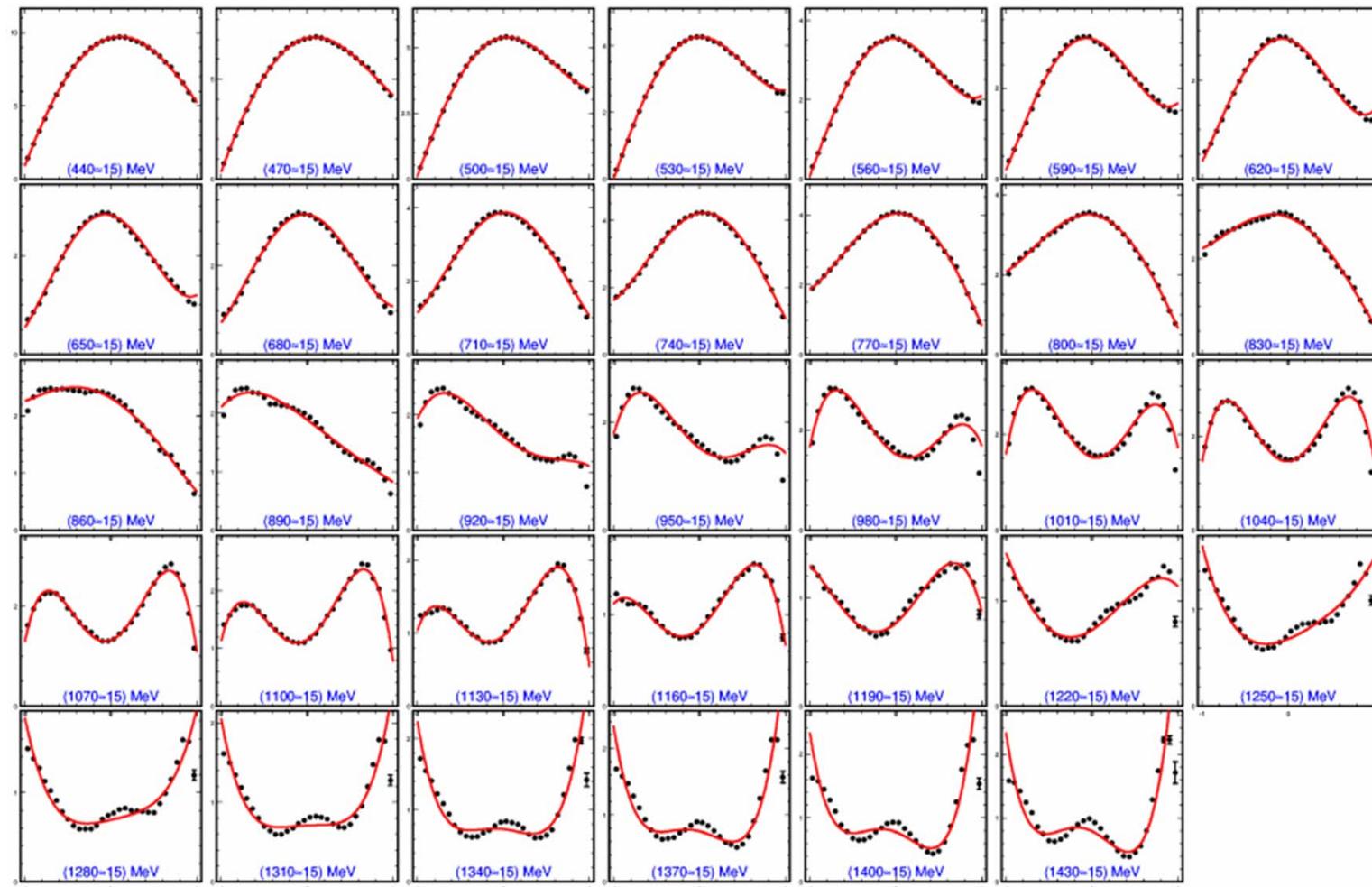


S+P waves are only good up to about 400 MeV

# Legendre expansion of the diff. cross section



$$\frac{d\sigma}{d\Omega} = \sum_{k=0}^{2\ell_{max}} A_k^\sigma(W) P_k(\cos \theta) \quad \ell_{max} = 2$$

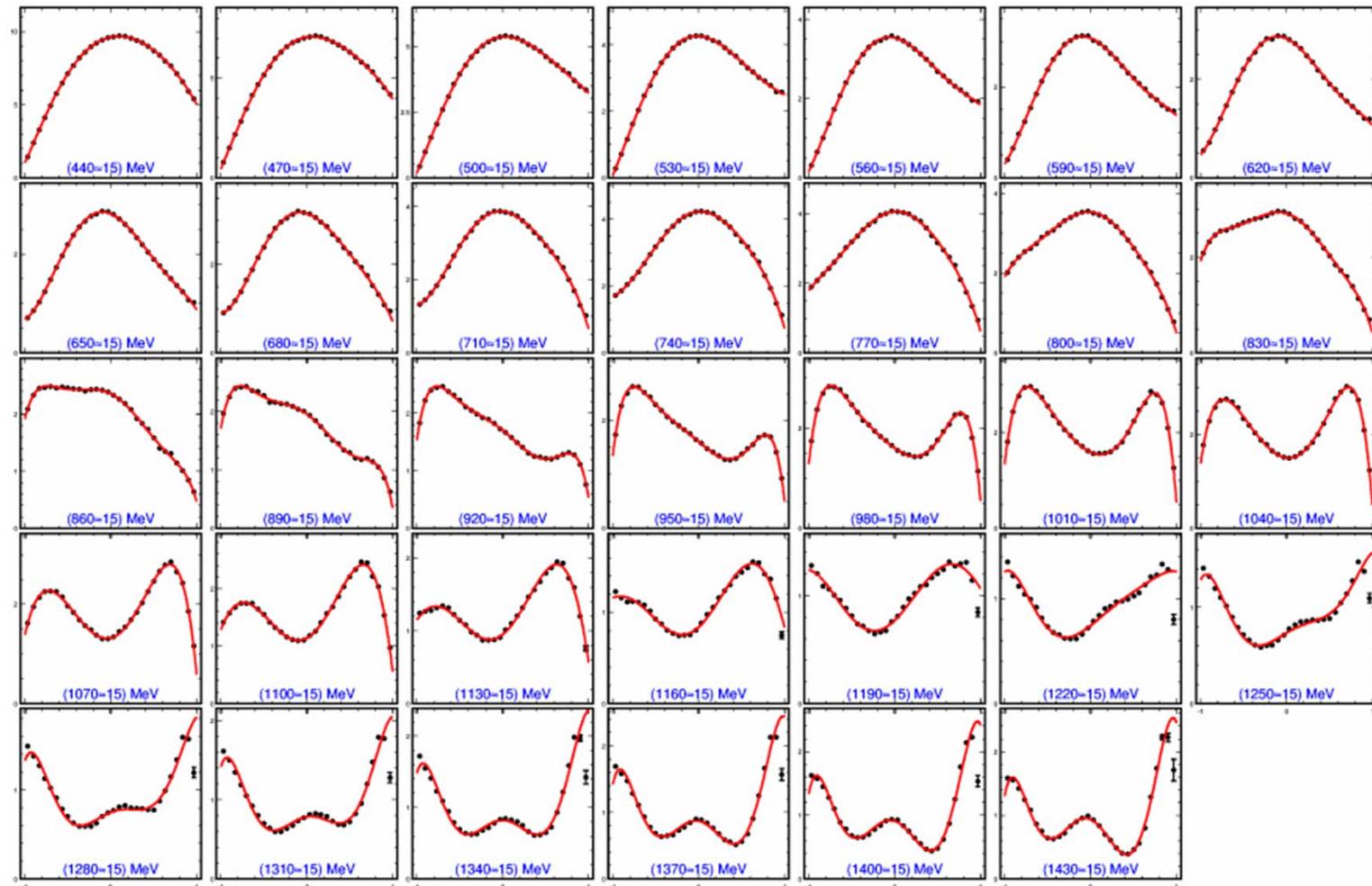


S+P+D waves are good up to about 850 MeV

# Legendre expansion of the diff. cross section



$$\frac{d\sigma}{d\Omega} = \sum_{k=0}^{2\ell_{max}} A_k^\sigma(W) P_k(\cos \theta) \quad \ell_{max} = 3$$

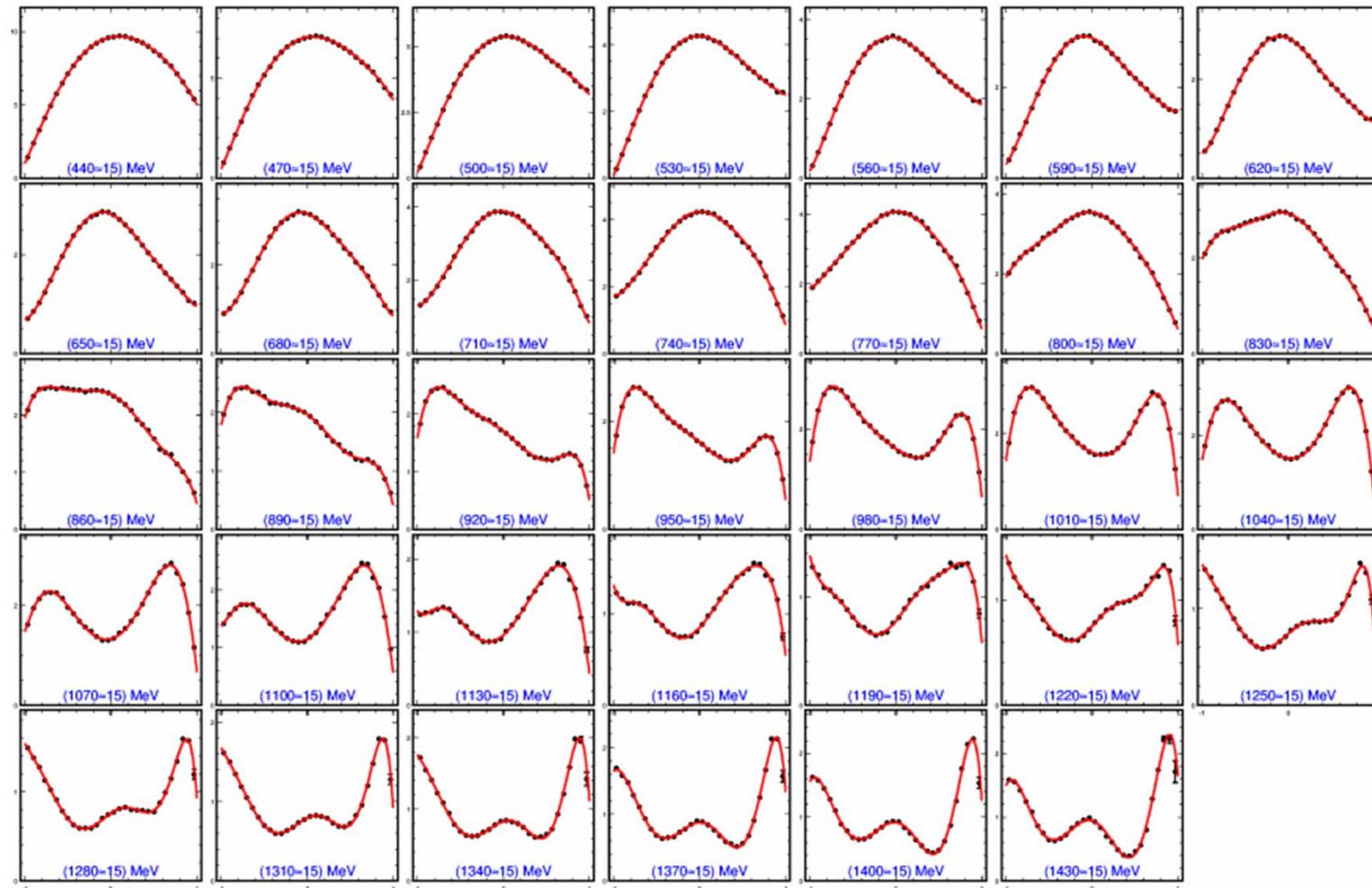


F waves become important around 1 GeV

# Legendre expansion of the diff. cross section



$$\frac{d\sigma}{d\Omega} = \sum_{k=0}^{2\ell_{max}} A_k^\sigma(W) P_k(\cos \theta) \quad \ell_{max} = 4$$

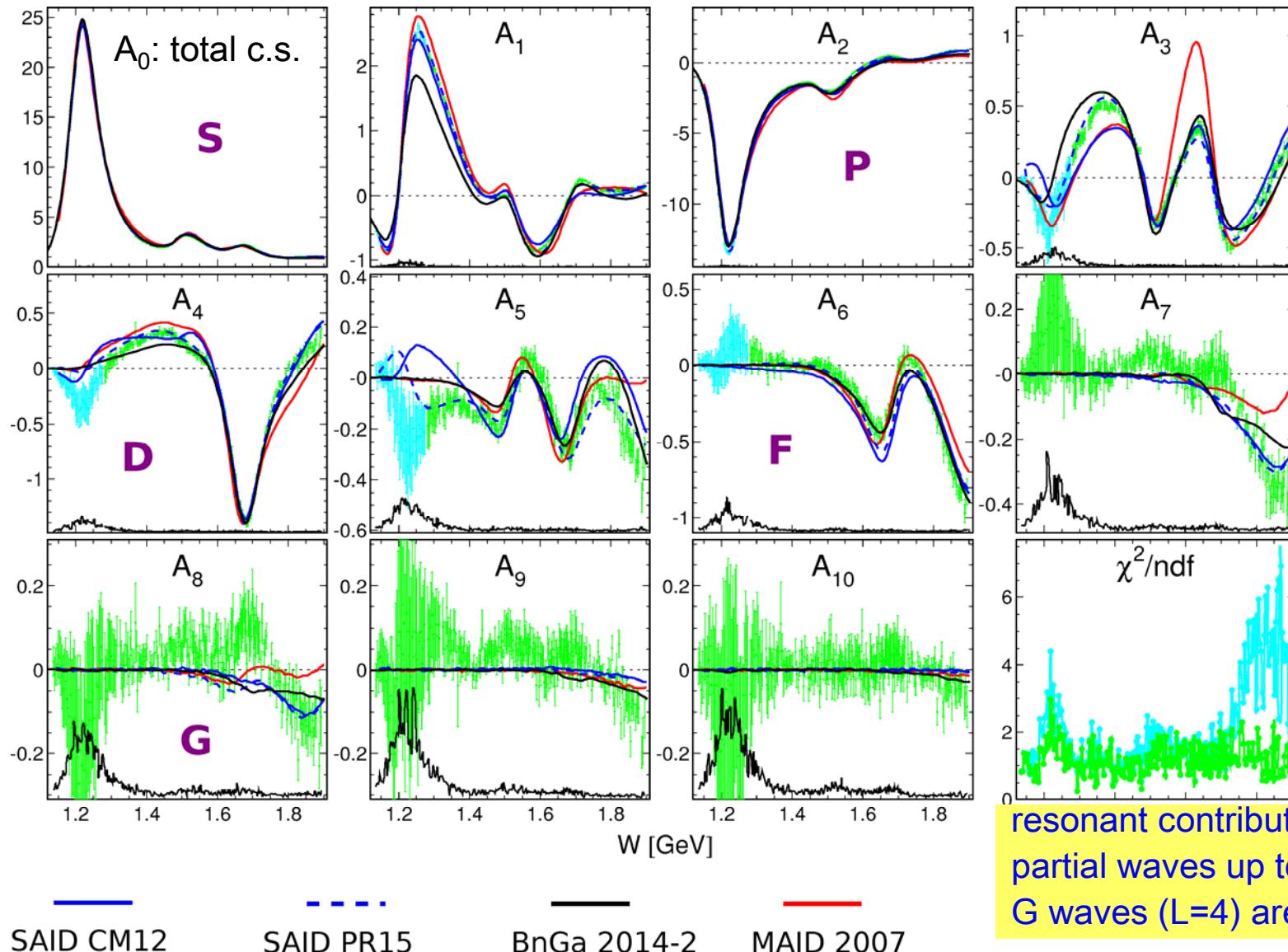


around 1.2 GeV also G waves become clearly visible in forward direction

# Legendre expansion of differential cross section



P. Adlarson et al (MAMI A2 Collab.) Phys. Rev. C92 (2015) 2, 024617



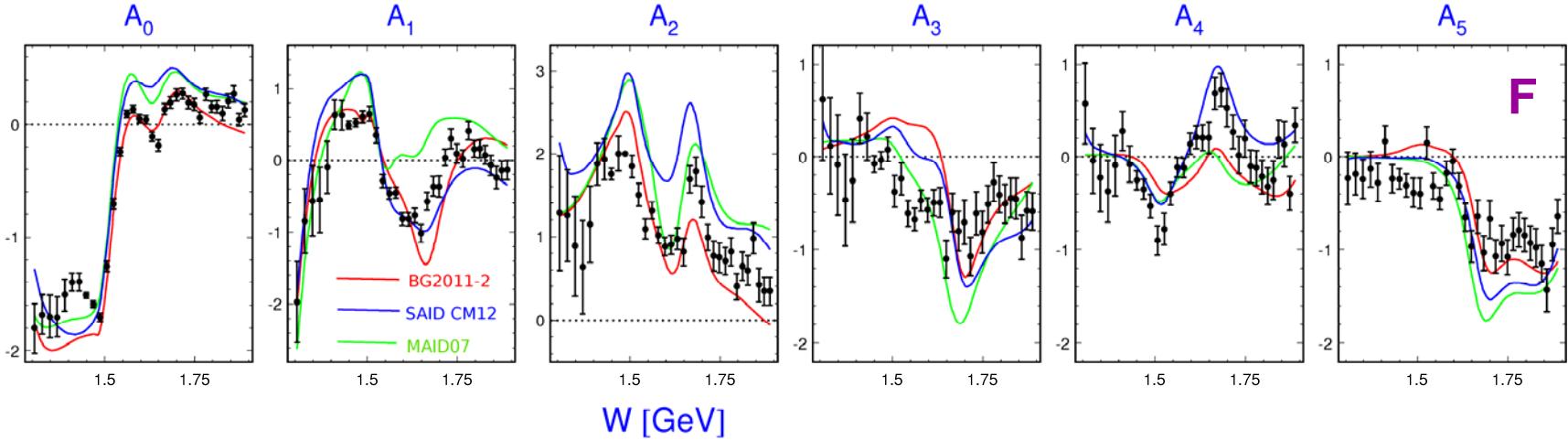
# Legendre expansion of polarization observables for $\gamma, \pi^0$

A2 Collaboration at MAMI, ready for publication (2015)



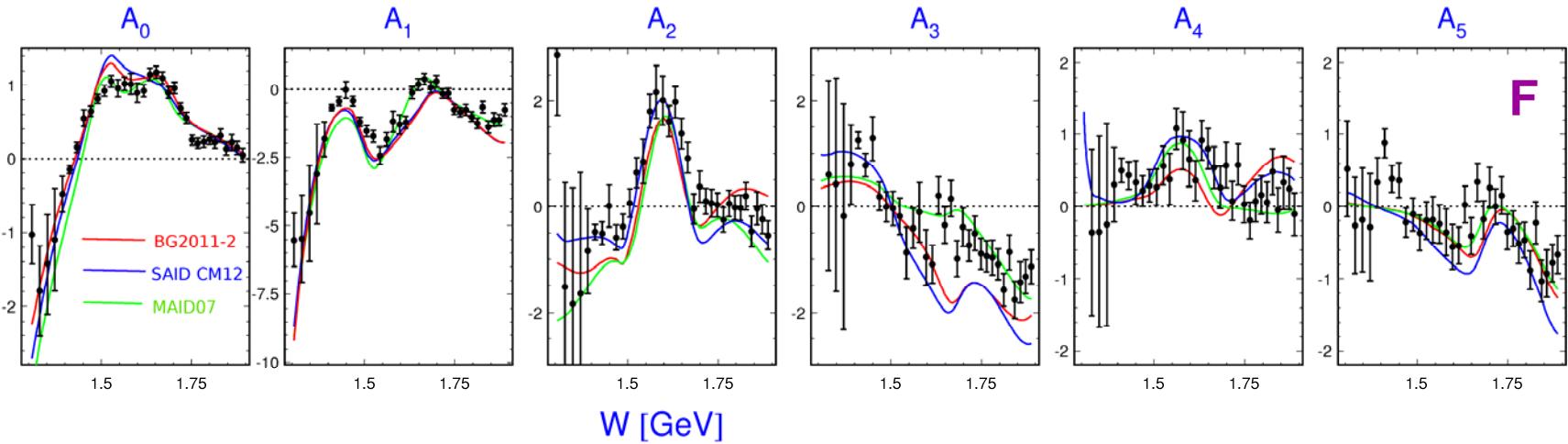
## T observable (target polarization)

$$T \cdot \frac{d\sigma}{d\Omega} = \sin\theta \sum_{n=0}^{2\ell_{max}} A_T^n P_n(\cos\theta)$$



## F observable (beam-target double polarization)

$$F \cdot \frac{d\sigma}{d\Omega} = \sin\theta \sum_{n=0}^{2\ell_{max}} A_F^n P_n(\cos\theta)$$

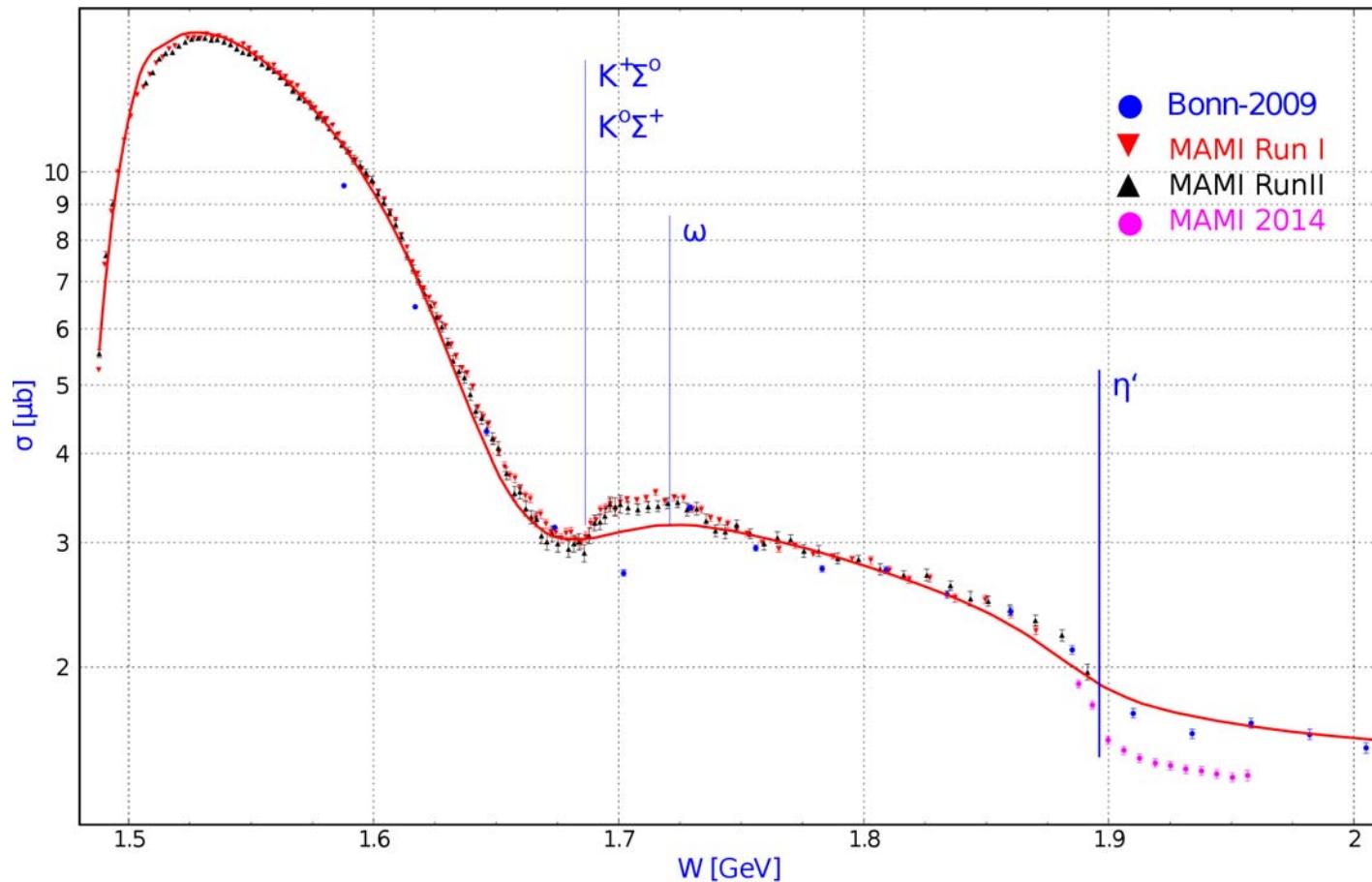


# $\eta$ Photoproduction at MAMI



$\gamma p \rightarrow \eta p$

MAMI A2 Collab, preliminary

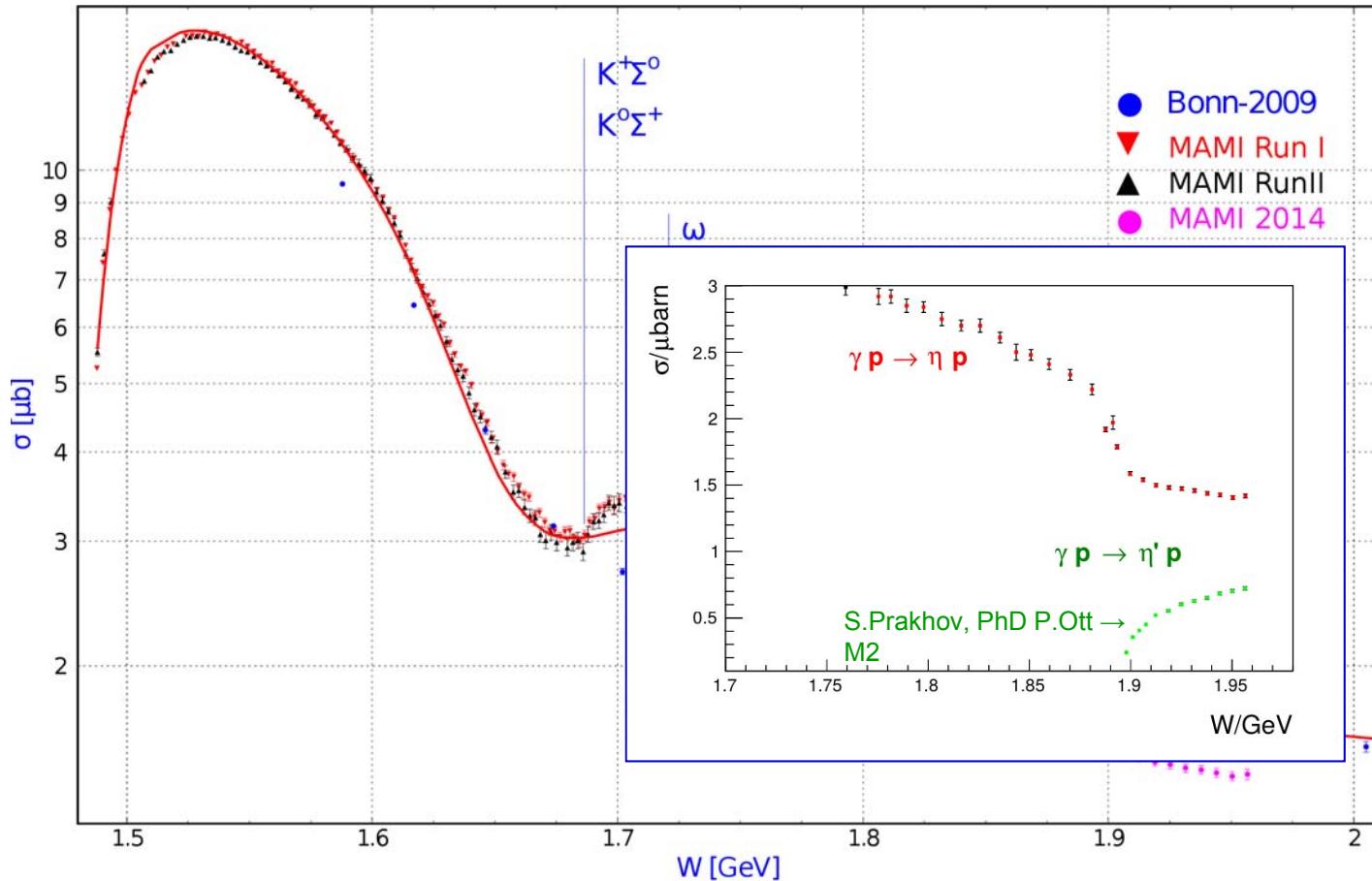


# $\eta$ Photoproduction at MAMI

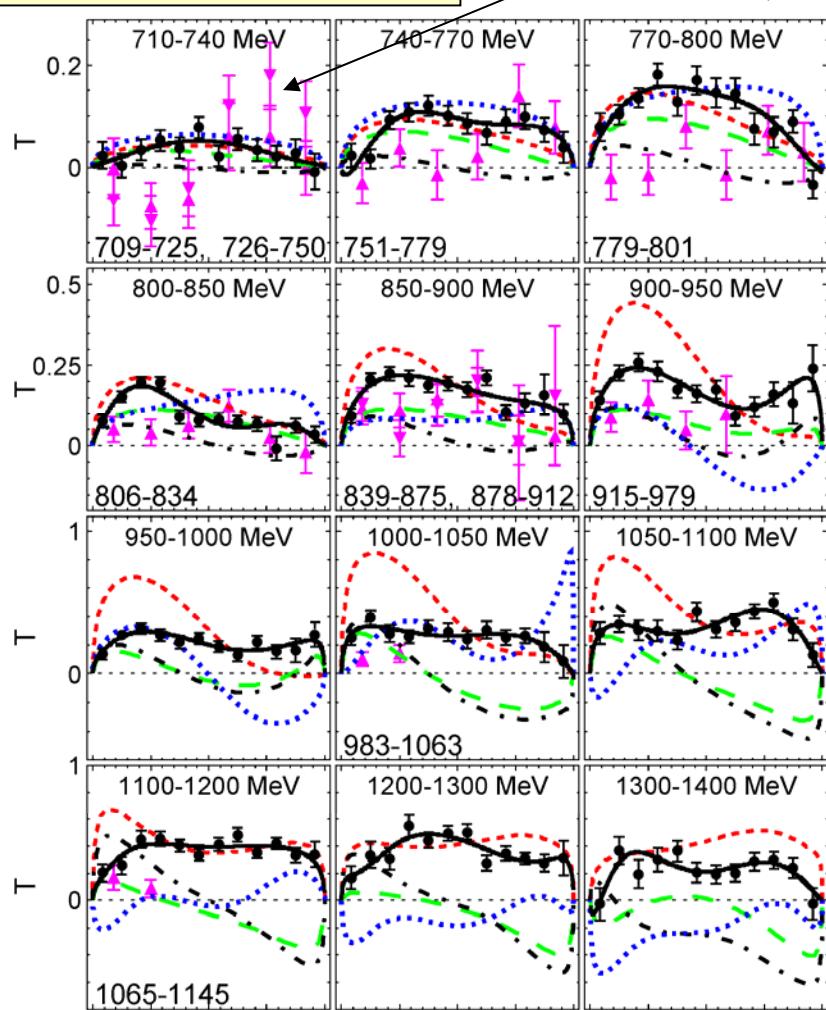


$\gamma p \rightarrow \eta p$

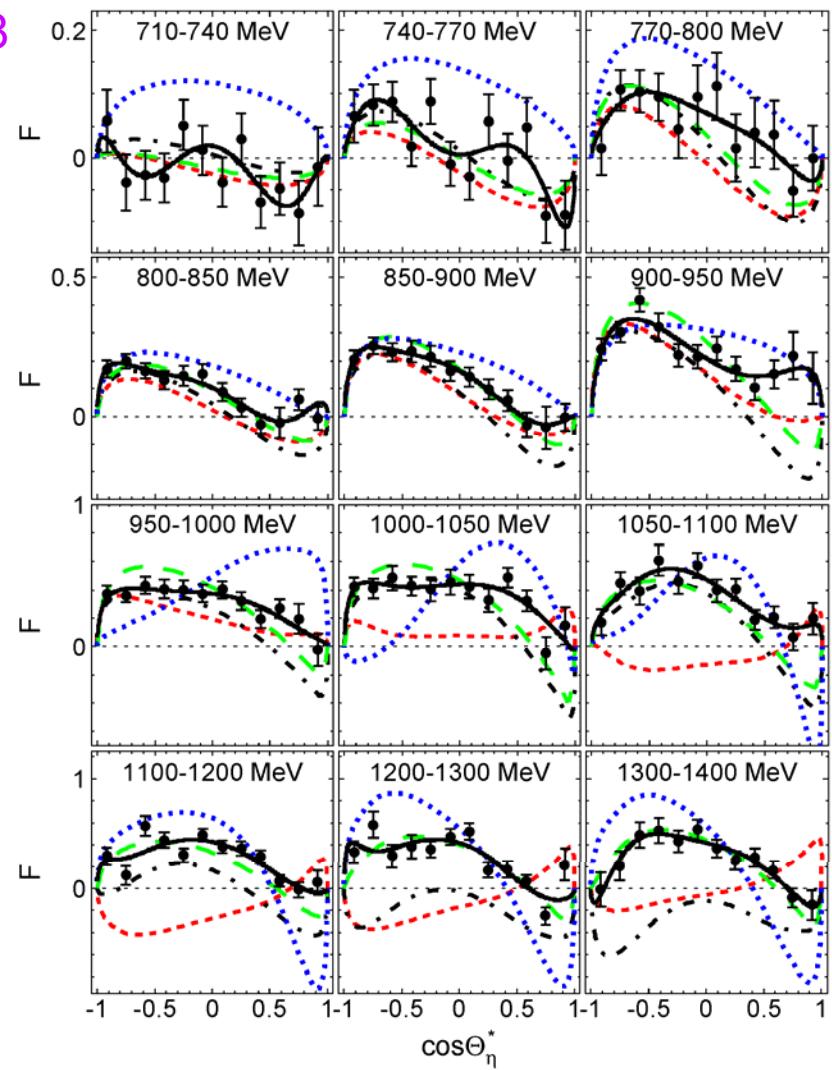
MAMI A2 Collab, preliminary



# Target Polarization and Beam-Target Asymmetry for $\gamma, \eta$



A2 Collaboration at MAMI  
Phys. Rev. Lett. 113 (2014)





# The MAID Ansatz

$$t_{\gamma\pi}(W, Q^2) = t_{\gamma\pi}^B(W, Q^2) + t_{\gamma\pi}^R(W, Q^2)$$

total amplitude

$$t_{\gamma\pi}^{B,\alpha}(W, Q^2) = v_{\gamma\pi}^{B,\alpha}(W, Q^2) [1 + it_{\pi N}^\alpha(W)]$$

unitarized  
background amplitude

$$t_{\gamma\pi}^{R,\alpha}(W, Q^2) = \bar{\mathcal{A}}_\alpha^R(W, Q^2) \frac{f_{\gamma N}(W) \Gamma_{tot}(W) M_R f_{\pi N}(W)}{M_R^2 - W^2 - i M_R \Gamma_{tot}(W)} e^{i \phi_R(W, Q^2)}$$

unitarized  
reson. amplitude

energy-dep. widths

$$\Gamma_{\pi N}(W) = \beta_\pi \Gamma_R \left( \frac{q(W)}{q_R} \right)^{2l+1} \left( \frac{X_R^2 + q_R^2}{X_R^2 + q^2(W)} \right)^\ell \frac{M_R}{W}$$

$$f_{\pi N}(W) = C_{\pi N} \left[ \frac{1}{(2j+1)\pi} \frac{\kappa(W)}{q(W)} \frac{M_N}{M_R} \frac{\Gamma_{\pi N}(W)}{\Gamma_{tot}^2(W)} \right]^{1/2}$$

vertex functions  
depend only on W

$$f_{\gamma N}(W) = \left( \frac{\kappa(W)}{\kappa_R} \right)^n \left( \frac{X_R^2 + \kappa_R^2}{X_R^2 + \kappa^2(W)} \right)$$

phenomenological parametrization of transition form factors:

$$\bar{\mathcal{A}}_\alpha^R(W, Q^2) \approx \bar{\mathcal{A}}_\alpha(Q^2) = \bar{\mathcal{A}}_\alpha(0)(1 + a_1 Q^2 + a_2 Q^4 + a_3 Q^6 + a_4 Q^8) e^{-b_1 Q^2}$$



# The MAID Ansatz without unitarity, e.g. for $\gamma, \eta$

$$t_{\gamma\pi}(W, Q^2) = t_{\gamma\pi}^B(W, Q^2) + t_{\gamma\pi}^R(W, Q^2)$$

total amplitude

$$t_{\gamma\pi}^{B,\alpha}(W, Q^2) = v_{\gamma\pi}^{B,\alpha}(W, Q^2) [1 + \cancel{i\epsilon_{\pi N}^\alpha(W)}]$$

unitarized  
background amplitude

$$t_{\gamma\pi}^{R,\alpha}(W, Q^2) = \bar{\mathcal{A}}_\alpha^R(W, Q^2) \frac{f_{\gamma N}(W) \Gamma_{tot}(W) M_R f_{\pi N}(W)}{M_R^2 - W^2 - i M_R \Gamma_{tot}(W)} e^{\cancel{i\phi_R(W, Q^2)}}$$

unitarized  
reson. amplitude

$$\Gamma_{\pi N}(W) = \beta_\pi \Gamma_R \left( \frac{q(W)}{q_R} \right)^{2l+1} \left( \frac{X_R^2 + q_R^2}{X_R^2 + q^2(W)} \right)^\ell \frac{M_R}{W}$$

energy-dep. widths

$$f_{\pi N}(W) = C_{\pi N} \left[ \frac{1}{(2j+1)\pi} \frac{\kappa(W)}{q(W)} \frac{M_N}{M_R} \frac{\Gamma_{\pi N}(W)}{\Gamma_{tot}^2(W)} \right]^{1/2}$$

vertex functions  
depend only on W

$$f_{\gamma N}(W) = \left( \frac{\kappa(W)}{\kappa_R} \right)^n \left( \frac{X_R^2 + \kappa_R^2}{X_R^2 + \kappa^2(W)} \right)$$

phenomenological parametrization of transition form factors:

$$\bar{\mathcal{A}}_\alpha^R(W, Q^2) \approx \bar{\mathcal{A}}_\alpha(Q^2) = \bar{\mathcal{A}}_\alpha(0)(1 + a_1 Q^2 + a_2 Q^4 + a_3 Q^6 + a_4 Q^8) e^{-b_1 Q^2}$$

# EtaMAID update for coupled $\gamma,\eta$ and $\gamma,\eta'$



(with V. Kashevarov, Mainz)

*still very preliminary, 28.09.2015*

## Resonances in $\eta$ MAID-2003

$D_{13}(1520)****$   
 $F_{15}(1680)****$

$S_{11}(1535)****$   
 $D_{13}(1700)***$

$S_{11}(1650)****$   
 $P_{11}(1710)***$

$D_{15}(1675)****$   
 $P_{13}(1720)****$

## Additional resonances in $\eta$ MAID-2015d

$F_{15}(1860)**$   
 $P_{13}(1900)***$   
 $D_{13}(2150)**$   
 $P_{11}(1440)****$

$D_{13}(1875)***$   
 $F_{17}(1990)**$   
 $G_{17}(2190)****$   
 $P_{11}(2300)**$

$P_{11}(1880)**$   
 $F_{15}(2000)**$   
 $H_{19}(2220)****$   
 $D_{15}(2570)**$

$S_{11}(1895)**$   
 $D_{15}(2060)**$   
 $G_{19}(2250)****$

red marked resonances have very small contributions and were excluded from fit

# EtaMAID update for coupled $\gamma,\eta$ and $\gamma,\eta'$



(with V. Kashevarov, Mainz)

*data used in the fit*

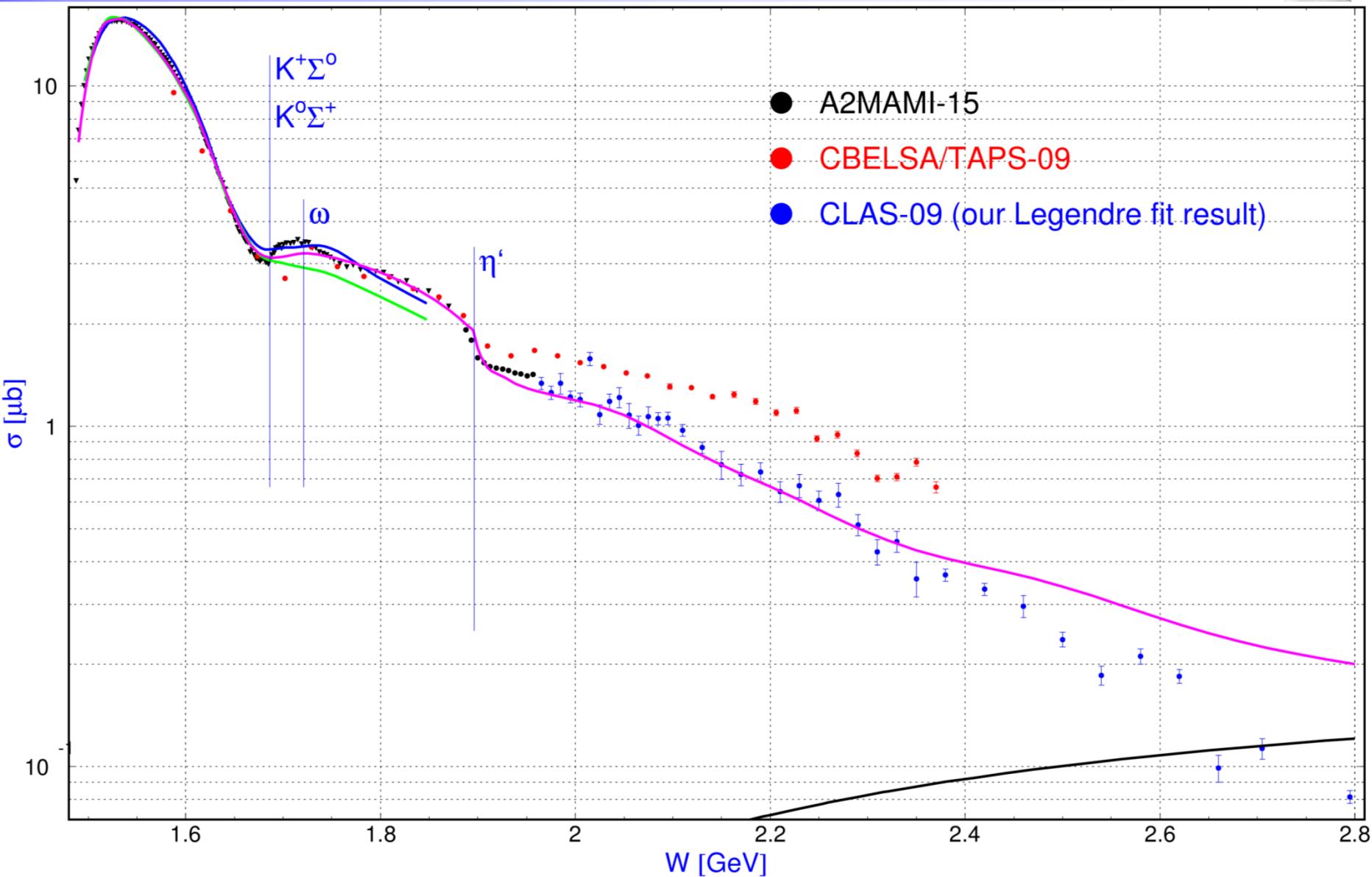
$\gamma p \rightarrow \eta p$

• $d\sigma/d\Omega$	W=1.488 – 1.957 GeV	MAMI 2015, Prakhov, preliminary
• $d\sigma/d\Omega$	W=1.965 – 2.075 GeV	CLAS 2009, PR C80 (2009) 045213
• T	W=1.497 – 1.848 GeV	MAMI 2014, PRL 113 (2014) 102001
• F	W=1.497 – 1.848 GeV	MAMI 2014, PRL 113 (2014) 102001
• $\Sigma$	W=1.496 – 1.908 GeV	GRAAL 2007, EPJ A33 (2007) 169
• E	W=1.525 – 2.125 GeV	CLAS 2015, arXiv:1507.00325v1

$\gamma p \rightarrow \eta' p$

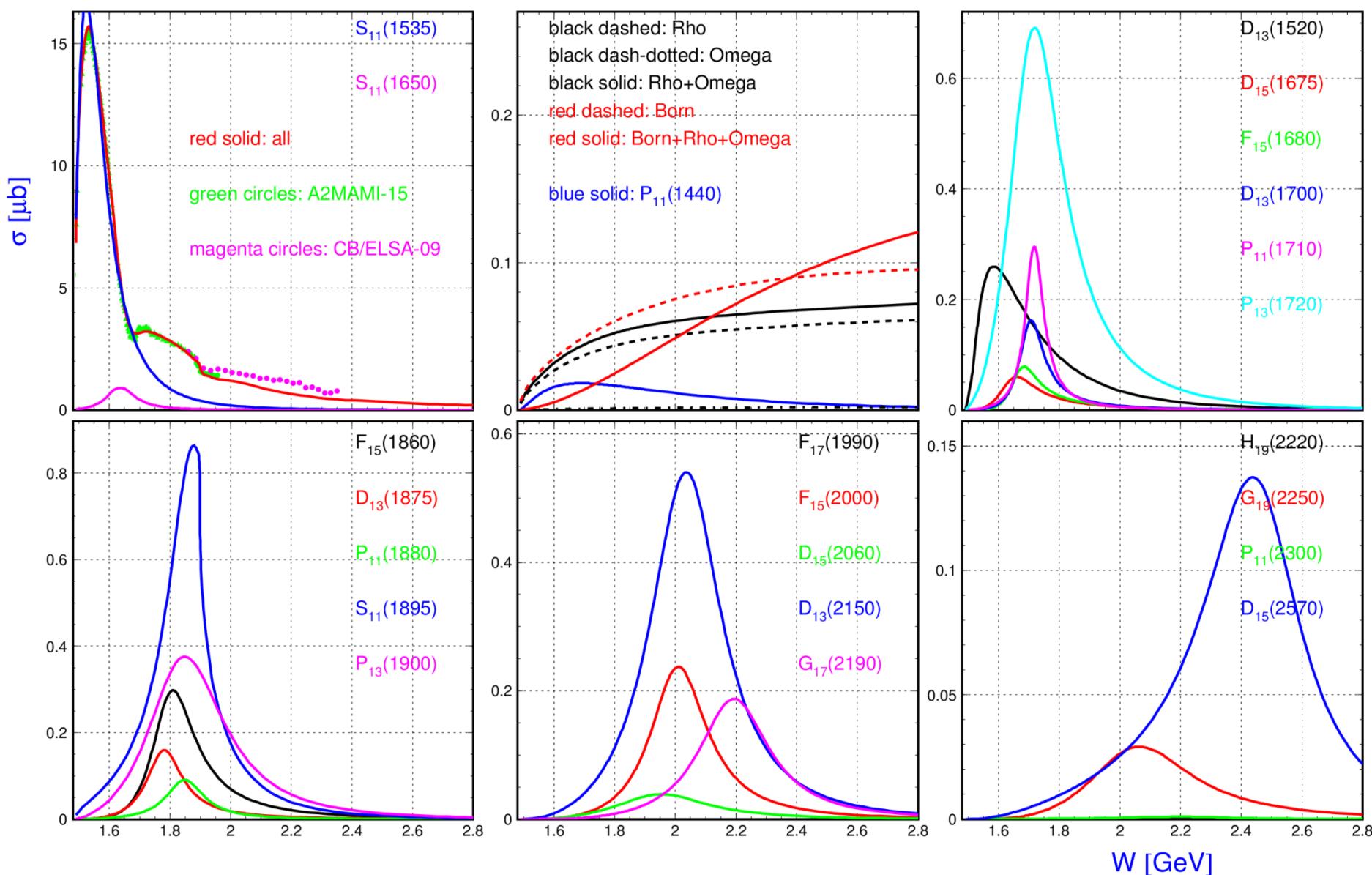
• $d\sigma/d\Omega$	W=1.898 – 1.956 GeV	MAMI 2015, Prakhov, preliminary
• $d\sigma/d\Omega$	W=1.925 – 2.045 GeV	CLAS 2009, PR C80 (2009) 045213
• $\Sigma$	W=1.903 – 1.913 GeV	GRAAL 2015, EPJ A51 (2015) 77

overall  $\chi^2 \sim 3.8$ , below  $\eta'$  threshold  $\sim 3.4$

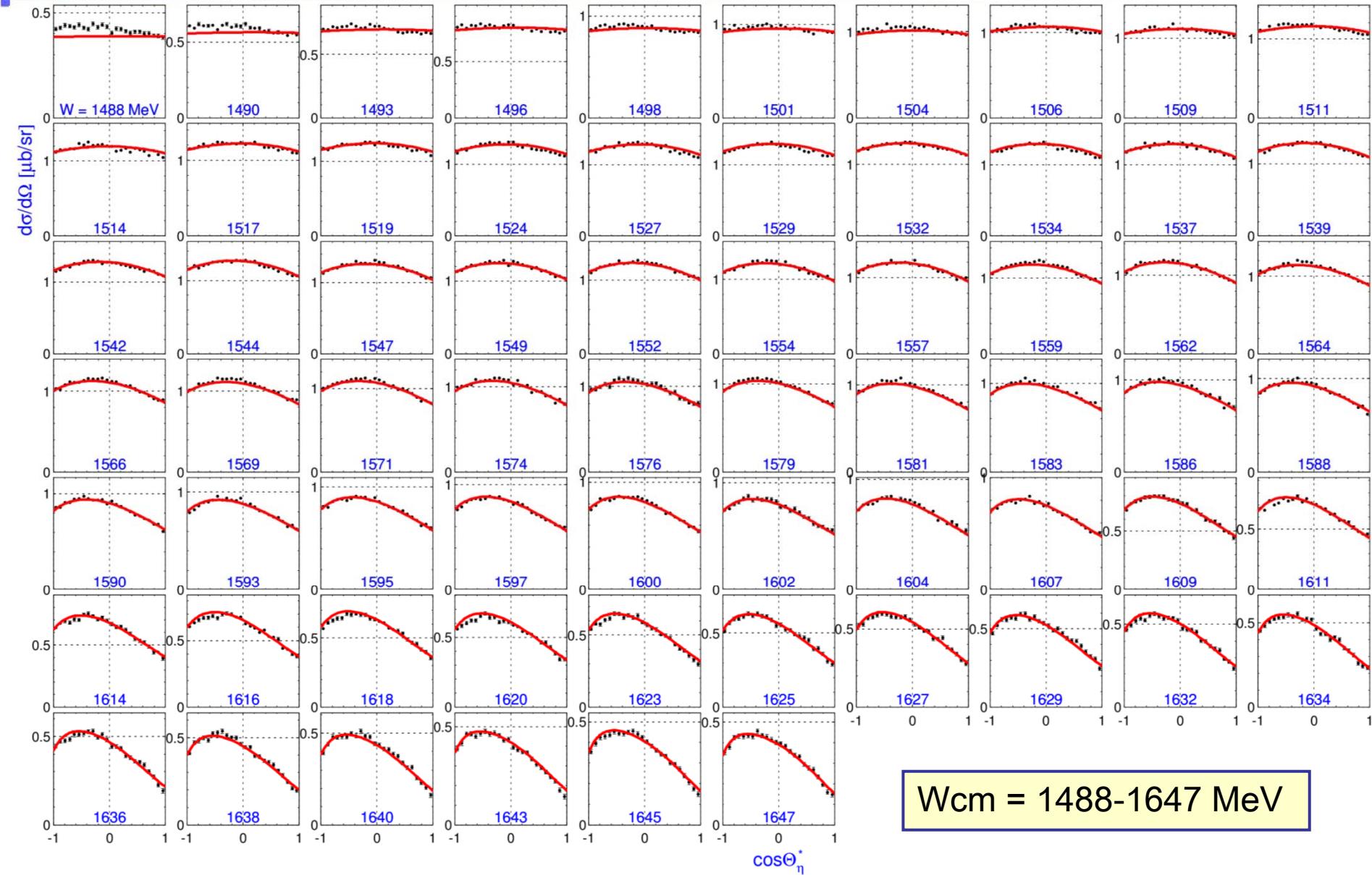


Blue line:  $\eta$ MAID-2003  
Green line:  $\eta$ MAID-2003regge

magenta line:  $\eta$ MAID-2015d  
black line:  $\eta$ MAID-2015d background



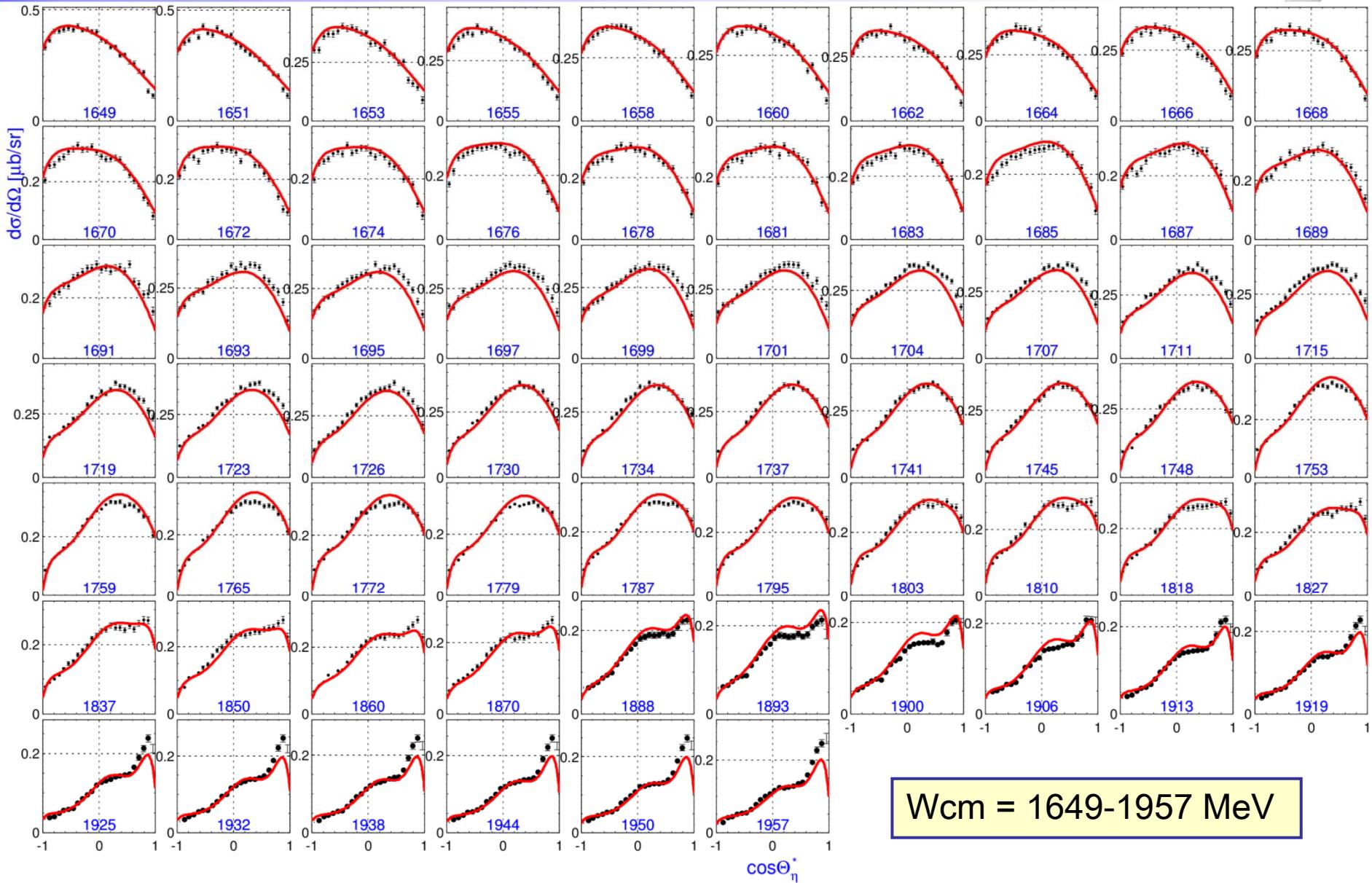
dominant states: S<sub>11</sub>(1535), S<sub>11</sub>(1650), S<sub>11</sub>(1895), P<sub>13</sub>(1720), D<sub>13</sub>(2120)<sup>new</sup>, P<sub>13</sub>(1900), ...

$\gamma p \rightarrow \eta p$  $\eta MAID-2015d: differential cross sections, Mainz data$ 

black circles: A2MAMI-15

red: ηMAID-2015d

$\gamma p \rightarrow \eta p$   *$\eta$ MAID-2015d: differential cross sections, Mainz data*

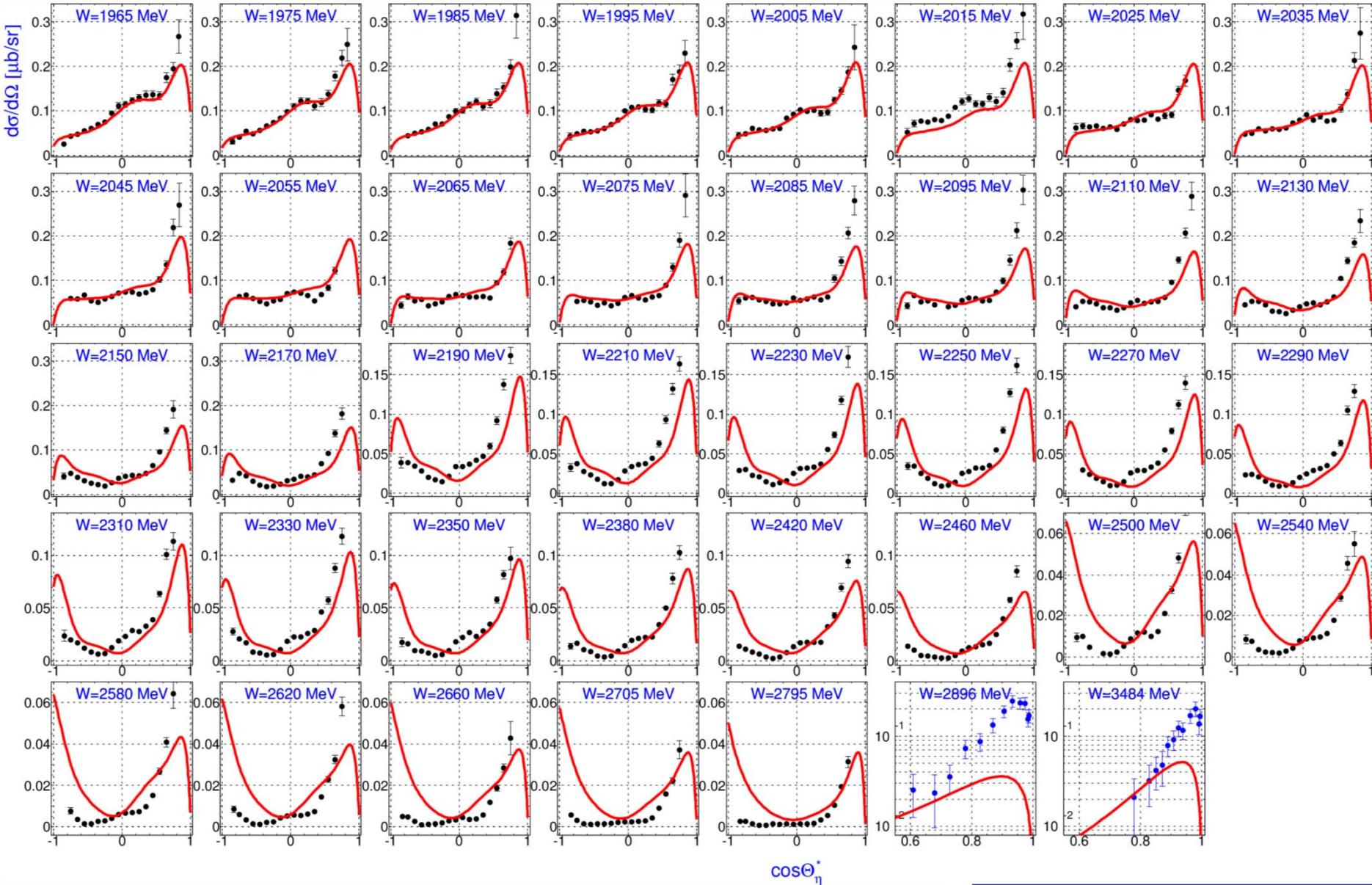


$W_{cm} = 1649-1957$  MeV

black circles: A2MAMI-15

red:  $\eta$ MAID-2015d

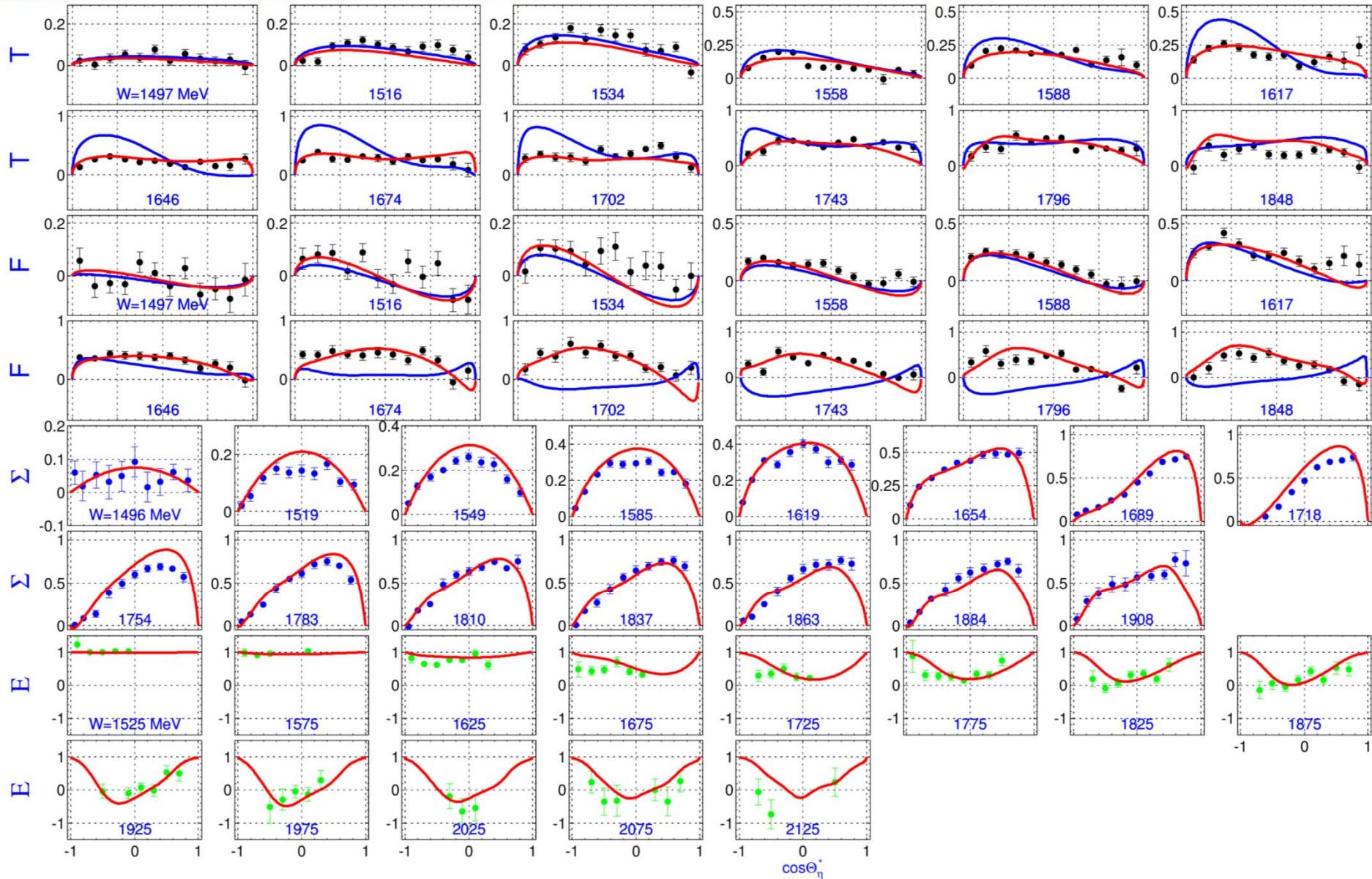
$\gamma$  p  $\rightarrow$   $\eta$  p    $\eta$ MAID-2015d: differential cross sections JLab/DESY



black circles: CLAS-09 (in the fit were included data up to  $W=2075$  MeV)  
 Blue circles: DESY-70  
 red line:  $\eta$ MAID-2015d

Wcm = 1965-3484 MeV

$\gamma p \rightarrow \eta p$   $\eta$ MAID-2015d:  $T, F, \Sigma, E$  data from Mainz/GRAAL/JLab

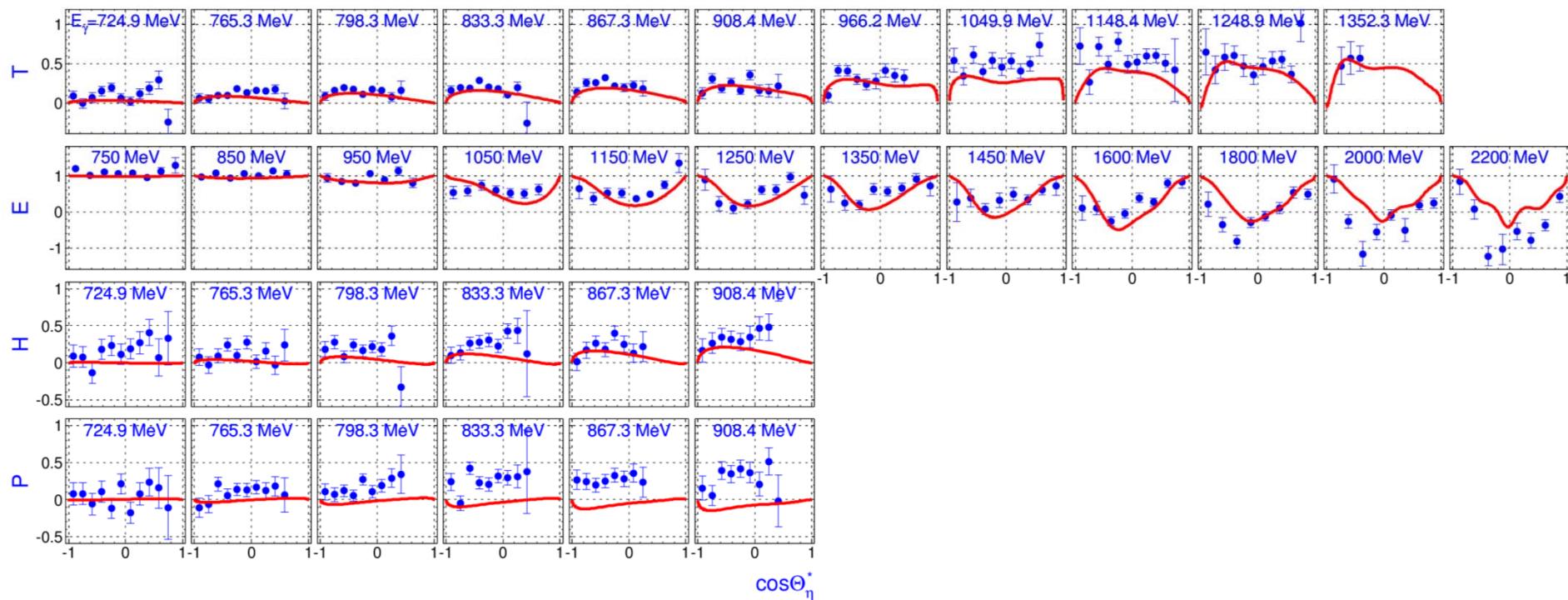
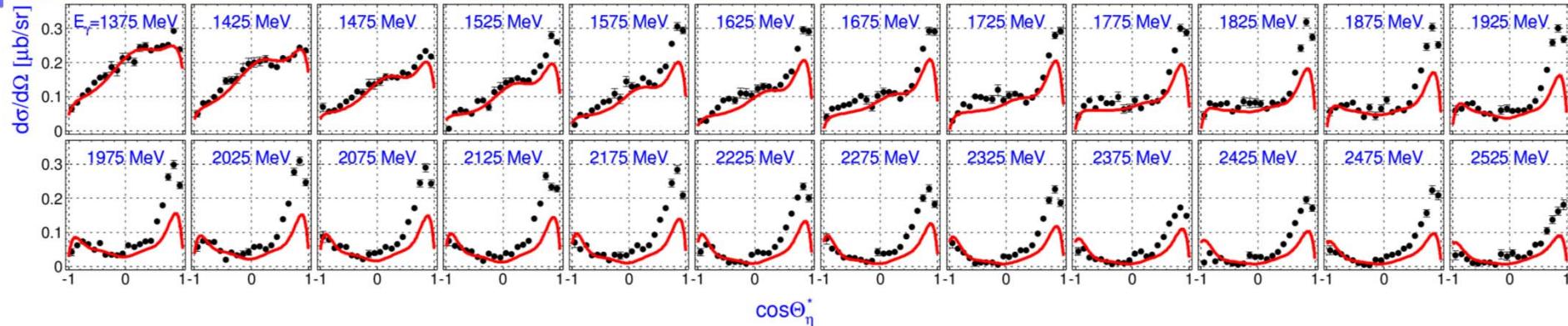


black circles: A2MAMI-14

blue circles: GRAAL-07; green: CLAS-15

blue lines:  $\eta$ MAID-2003

red lines:  $\eta$ MAID-2015d

$\gamma p \rightarrow \eta p$  $\eta\text{MAID-2015d}$  predictions for CBELSA/TAPS data

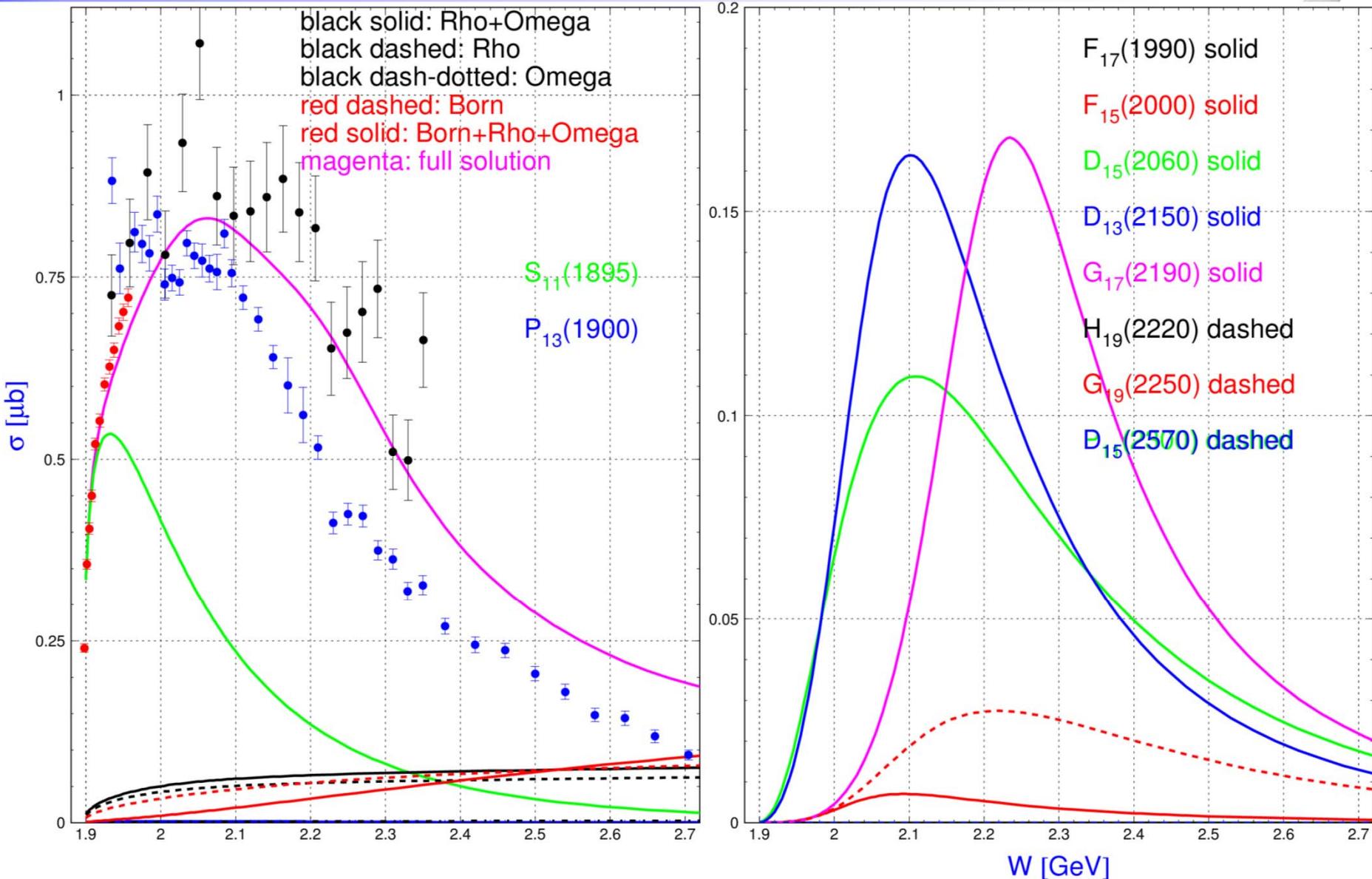
black circles: CBELSA/TAPS-09,

blue circles: CBELSA/TAPS-15 (J. Hartmann: T,P,H,

red lines: ηMAID-2015d

J. Müller: E, preliminary)

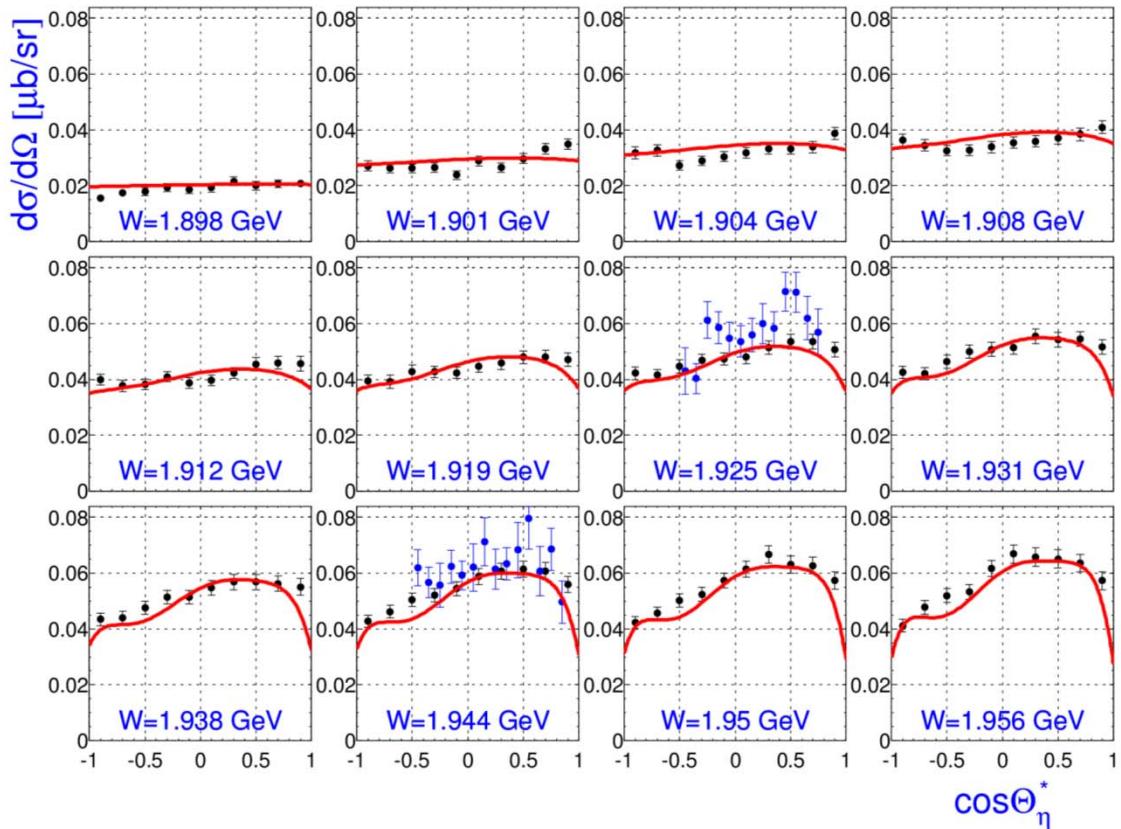
$\gamma$  p →  $\eta'$  p   *ηMAID-2015d: total cross section Mainz/Bonn/JLab data*



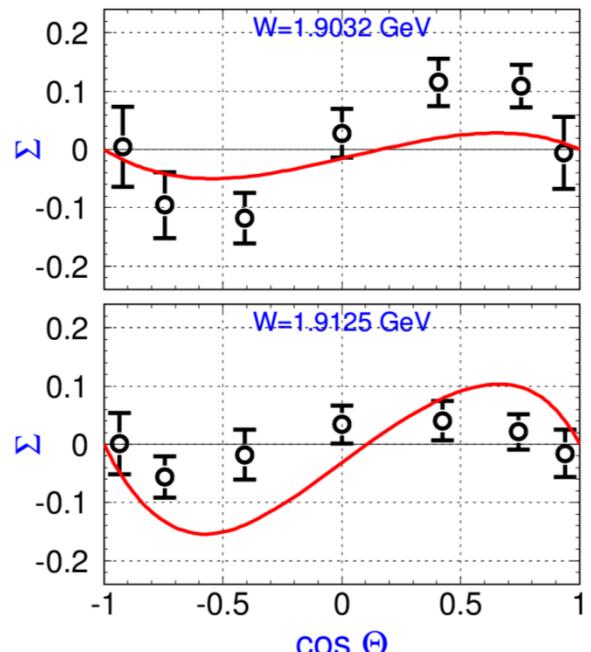
Red circles: A2MAMI-15,    black circles: CBELSA/TAPS-09,    blue circles: CLAS-09 from Legendre fit



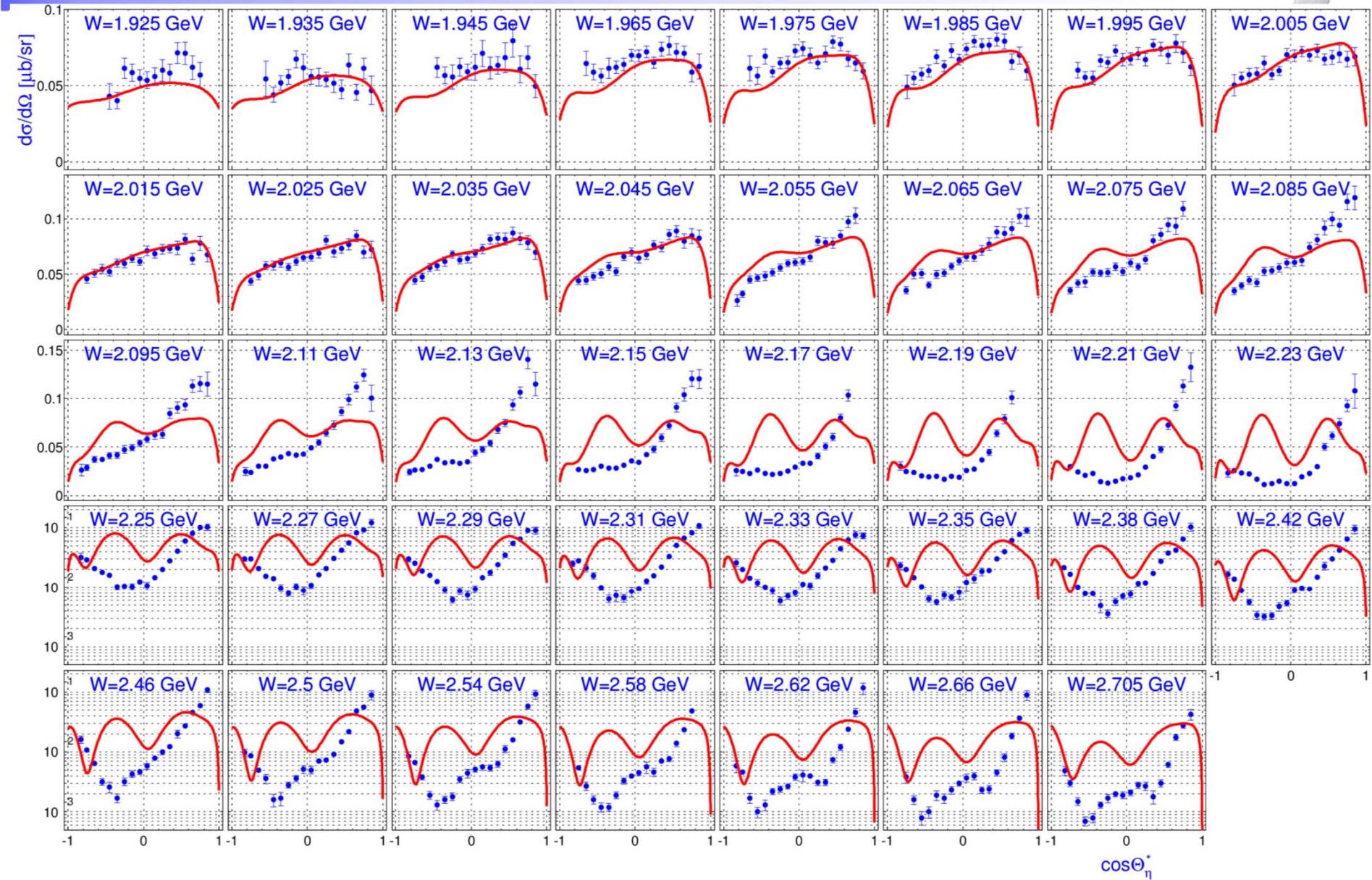
MAMI - A2 data 2015, preliminary



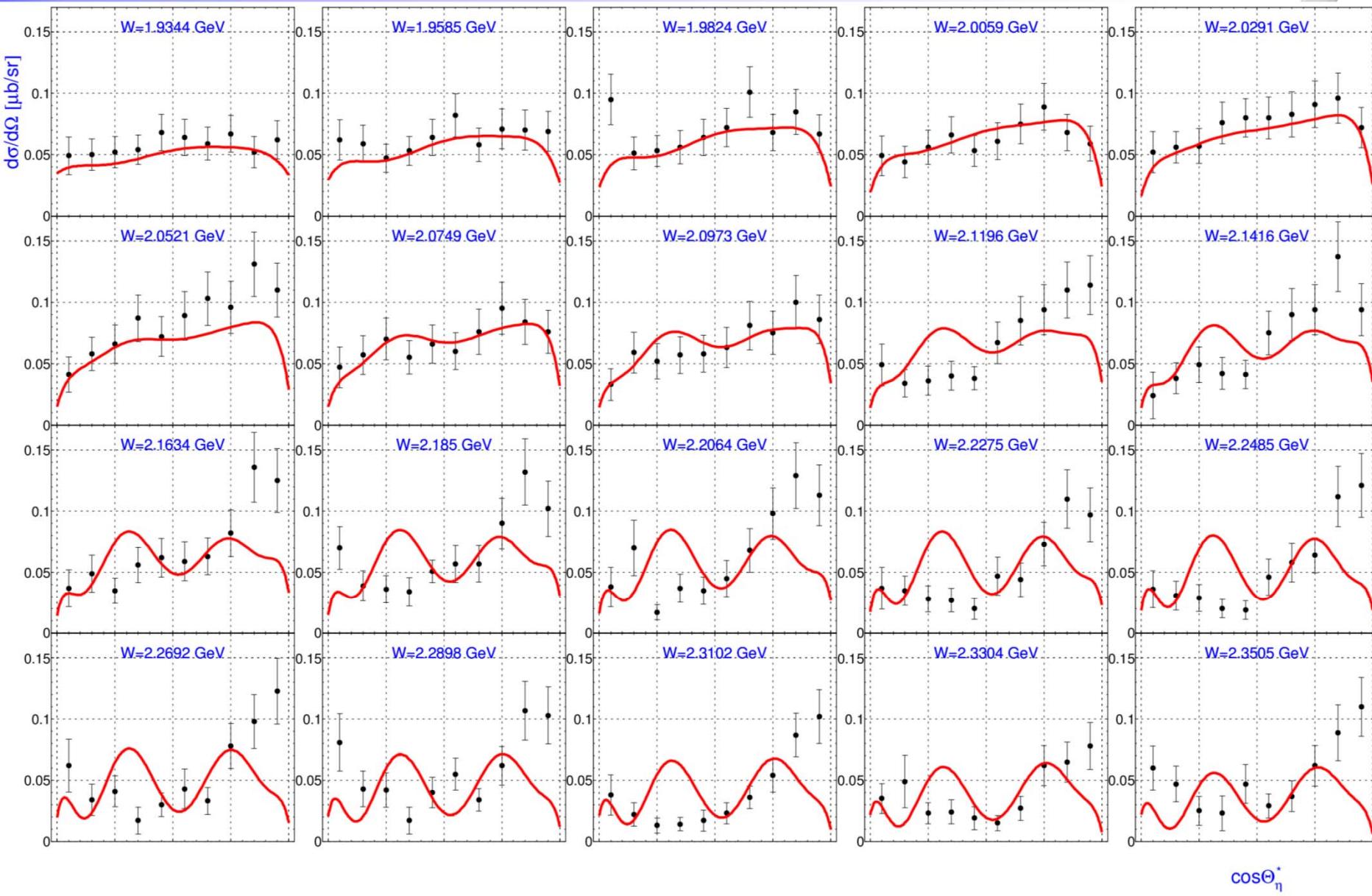
GRAAL data 2015



closed circles: A2MAMI-15; open circles : GRAAL-15; Red lines:  $\eta MAID-2015d$

$\gamma p \rightarrow \eta' p$  $\eta MAID-2015d$  and  $CLAS-09$  data

blue circles: CLAS-09 (in the fit were included data up to  $W=2045$  MeV) red line:  $\eta MAID-2015d$

$\gamma p \rightarrow \eta' p$  $\eta MAID-2015d$  and CBELSA/TAPS-09 data

black circles: CBELSA/TAPS-09 (not included in the fit)

red line:  $\eta$  MAID-2015d

# Regge parametrization of $t$ - channel vector meson exchange



around the  $\eta'$  threshold, the onset of the Regge regime becomes visible at forward angles

for  $\eta'$  photoproduction one can use Regge from threshold on in a single-channel analysis

for  $\eta$  photoproduction a smooth transition from poles to Regge has to be found

we plan to describe  $\eta$  and  $\eta'$  photoproduction in a coupled-channel approach for the vector meson exchange we want to use the method of V. Mathieu, I. Danilkin et al, arXiv:1506.01764, for matching low-energy analysis with high-energy Regge parametr.

# Eta-MAID update 2015 with new resonances



Particle	$J^P$	overall	$N\gamma$	$N\pi$	$N\eta$	$N\sigma$	$N\omega$	$\Lambda K$	$\Sigma K$	$N\rho$	$\Delta\pi$
$N(1440)$	$1/2^+$	****	****	****	○	***			*	***	
$N(1520)$	$3/2^-$	****	****	****	***				***	***	
$N(1535)$	$1/2^-$	****	****	****	****				**	*	
$N(1650)$	$1/2^-$	****	****	****	***			***	**	**	***
$N(1675)$	$5/2^-$	****	****	****	*			*	*	***	
$N(1680)$	$5/2^+$	****	****	****	*	**			***	***	
$N(1700)$	$3/2^-$	***	**	***	*			*	*	*	***
$N(1710)$	$1/2^+$	****	****	****	***	**	****	***	*	**	
$N(1720)$	$3/2^+$	****	****	****	***			**	**	**	*
$N(1860)$	$5/2^+$	**		**	○				*	*	
$N(1875)$	$3/2^-$	***	***	*	○			**	***	**	***
$N(1880)$	$1/2^+$	**	*	*	○	**			*		
$N(1895)$	$1/2^-$	**	**	*	**				**	*	
$N(1900)$	$3/2^+$	***	***	**	**	**	***	**	*	**	
$N(1990)$	$7/2^+$	**	**	**					*		
$N(2000)$	$5/2^+$	**	**	*	**			**	*	**	
$N(2040)$	$3/2^+$	*		*	○						
$N(2060)$	$5/2^-$	**	**	**	*				**		
$N(2100)$	$1/2^+$	*		*	○						
$N(2120)$	$3/2^-$	**	**	**	○			*	*		
$N(2190)$	$7/2^-$	****	***	****	○			*	**	*	
$N(2220)$	$9/2^+$	****		****							
$N(2250)$	$9/2^-$	****		****	○						
$N(2300)$	$1/2^+$	**		**							
$N(2570)$	$5/2^-$	**		**	○						



7  $N^*$  in 2001/2003



15  $N^*$  new in 2015

only 3  $N^*$  resonances  
in PDG below 2.6 GeV,  
where we do not find  
evidence for  $\gamma, \eta$

but everything is still preliminary

# resonance parameters (below $\eta'$ threshold)



N*	M	$\Gamma$	Br Eta	A1/2	A3/2	PDG stars for N $\eta$
D13(1520)	1513.	125.	0.11	-28.0	124.	***
S11(1535)	1534.	161.	42.	115.	---	****
S11(1650)	1645.	116.	-12.	45.	---	***
D15(1675)	1659.	165.	-1.9	11.	33.	*
F15(1680)	1680.	115.	0.10	-9.0	145.	*
D13(1700)	1711.	106.	-2.6	14.0	-37.	*
P11(1710)	1724.	80.	2.2	50.	---	***
<b>P13(1720)</b>	<b>1745.</b>	<b>268.</b>	<b>7.2</b>	<b>70.</b>	<b>30.</b>	<b>***</b>
F15(1860)	1819.	192.	1.1	-96.	-64.	new
D13(1875)	1831.	275.	-2.2	-59.	7.0	new
P11(1880)	1862.	158.	20.	-13.	---	new

# resonance parameters (above $\eta'$ threshold)



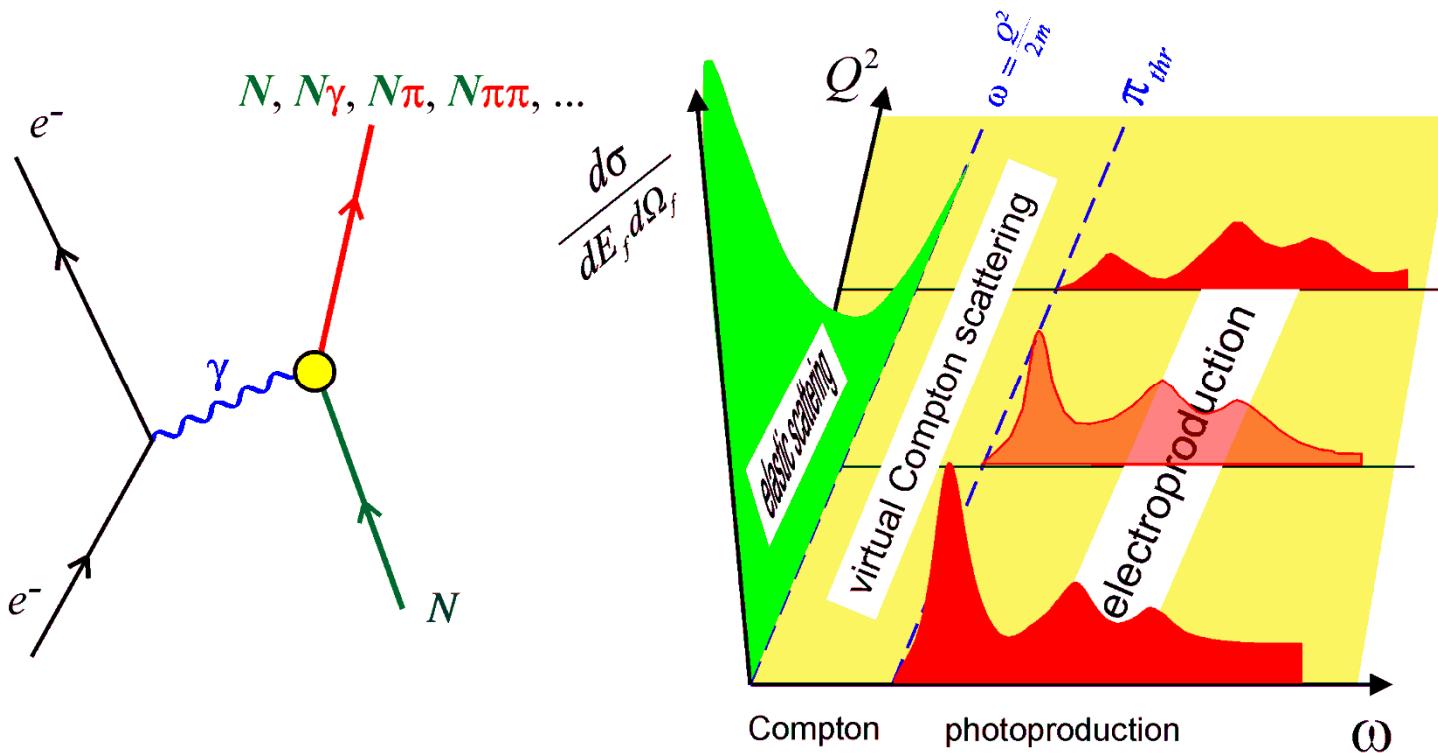
N*	M	$\Gamma$	Br Eta	Br Eta'	A1/2	A3/2	PDG stars for N $\eta$
S11(1895)	1895.	220.	26.	5.2	-40.	---	**
P13(1900)	1915.	383.	-14.	-0.017	44.	-21.	**
F17(1990)	1894.	114.	0.046	-0.00	-19.	-93.	new
F15(2000)	2029.	244.	-0.67	0.015	-141.	60.	**
D15(2060)	2014.	454.	0.42	-0.92	-102.	8.9	*
<b>D13(2120)</b>	<b>2064.</b>	<b>305.</b>	<b>1.17</b>	<b>-0.34</b>	<b>147.</b>	<b>-130.</b>	<b>new</b>
G17(2190)	2223.	303.	2.07	1.9	88.	-18.	new
G19(2250)	2149.	575.	-2.05	-2.0	46.	0.8	new
D15(2570)	2475.	398.	-4.08	-0.29	-20.	-63.	new



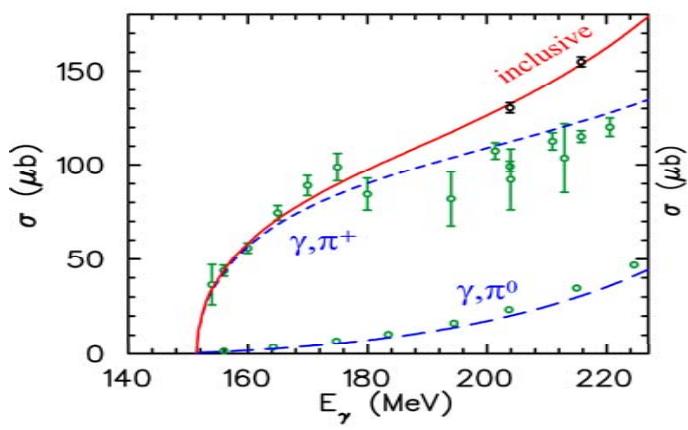
# 10 resonances according to their importance in $\sigma_{\text{total}}(\gamma, \eta)$

N*	M	$\Gamma$	Br Eta	A1/2	A3/2	PDG stars for Nn
S11(1535)	1534.	161.	42.	115.	---	**** 17 $\mu\text{b}$
S11(1650)	1645.	116.	-12.	45.	---	*** 1 $\mu\text{b}$
S11(1895)	1895.	220.	26. -40.		---	** 0.85 $\mu\text{b}$
<b>P13(1720)</b>	<b>1745.</b>	<b>268.</b>	<b>7.2</b>	<b>70.</b>	<b>30.</b>	<b>***</b> 0.70 $\mu\text{b}$
<b>D13(2120)</b>	<b>2064.</b>	<b>305.</b>	<b>1.17</b>	<b>147.</b>	<b>-130.</b>	<b>new</b> 0.55 $\mu\text{b}$
P13(1900)	1915.	383.	-14.	44.	-21.	** 0.40 $\mu\text{b}$
P11(1710)	1724.	80.	2.2	50.	---	*** 0.30 $\mu\text{b}$
F15(1860)	1819.	192.	1.1	-96.	-64.	new 0.30 $\mu\text{b}$
D13(1520)	1513.	125.	0.11	-28.0	124.	*** 0.25 $\mu\text{b}$
F15(2000)	2029.	244.	-0.67	-141.	60.	** 0.25 $\mu\text{b}$

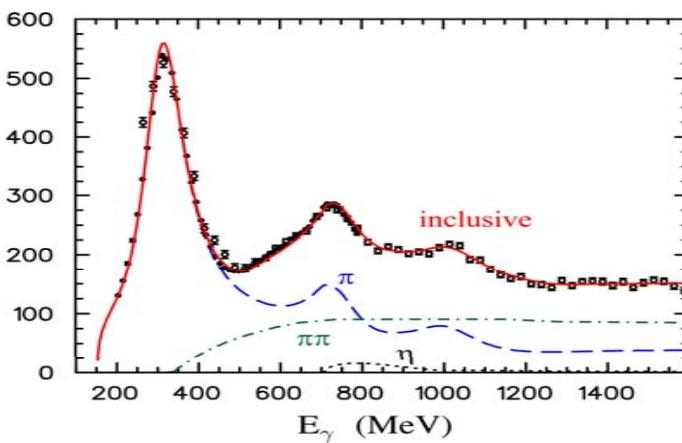
# from real to virtual photons



Threshold Region



Resonance Region



# The MAID Ansatz for $(e, e' \pi)$



unitarized resonance amplitude

$$t_{\gamma\pi}^{R,\alpha}(W, Q^2) = \bar{\mathcal{A}}_\alpha^R(W, Q^2) \frac{f_{\gamma N}(W) \Gamma_{tot}(W) M_R f_{\pi N}(W)}{M_R^2 - W^2 - i M_R \Gamma_{tot}(W)} e^{i\phi_R(W, Q^2)}$$

phenomenological parametrization of transition form factors:

$$\bar{\mathcal{A}}_\alpha^R(W, Q^2) \approx \bar{\mathcal{A}}_\alpha(Q^2) = \bar{\mathcal{A}}_\alpha(0)(1 + a_1 Q^2 + a_2 Q^4 + a_3 Q^6 + a_4 Q^8) e^{-b_1 Q^2}$$

last full version of  $\gamma, \pi$  and  $e, e' \pi$ : 2007

last transition form factors update: 2009



# MAID2007 was designed with minimal number of resonances

Particle	$J^P$	overall	$N\gamma$	$N\pi$	$N\eta$	$N\sigma$	$N\omega$	$\Lambda K$	$\Sigma K$	$N\rho$	$\Delta\pi$
$N(1440)$	$1/2^+$	****	****	*****		***			*	***	
$N(1520)$	$3/2^-$	****	****	*****	***				***	***	
$N(1535)$	$1/2^-$	****	****	*****	****				**	*	
$N(1650)$	$1/2^-$	****	****	*****	***			***	**	**	***
$N(1675)$	$5/2^-$	****	****	*****	*			*		*	***
$N(1680)$	$5/2^+$	****	****	*****	*	**			***	***	
$N(1700)$	$3/2^-$	***	**	***	*			*	*	*	***
$N(1710)$	$1/2^+$	****	****	****	***		**	*****	**	*	**
$N(1720)$	$3/2^+$	****	****	*****	***			**	**	**	*
$N(1860)$	$5/2^+$	**		**						*	*
$N(1875)$	$3/2^-$	***	***	*			**	***	**		***
$N(1880)$	$1/2^+$	**	*	*		**		*			
$N(1895)$	$1/2^-$	**	**	*	**			**	*		
$N(1900)$	$3/2^+$	***	***	**	**		**	***	**	*	**
$N(1990)$	$7/2^+$	**	**	**						*	
$N(2000)$	$5/2^+$	**	**	*	**			**	*	**	
$N(2040)$	$3/2^+$	*		*							
$N(2060)$	$5/2^-$	**	**	**	*				**		
$N(2100)$	$1/2^+$	*		*							
$N(2120)$	$3/2^-$	**	**	**				*	*		
$N(2190)$	$7/2^-$	****	***	****		*	**		*		
$N(2220)$	$9/2^+$	****		****							
$N(2250)$	$9/2^-$	****		****							
$N(2300)$	$1/2^+$	**		**							
$N(2570)$	$5/2^-$	**		**							

only  $7 N^*$  resonances  
in MAID2007



# MAID2007 was designed with minimal number of resonances

Particle	$J^P$	overall	$N\gamma$	$N\pi$	$N\eta$	$N\sigma$	$N\omega$	$\Lambda K$	$\Sigma K$	$N\rho$	$\Delta\pi$
$\Delta(1232)$	$3/2^+$	****	****	*****		F					
$\Delta(1600)$	$3/2^+$	***	***	***		O				*	***
$\Delta(1620)$	$1/2^-$	****	***	*****		r				***	***
$\Delta(1700)$	$3/2^-$	****	****	*****		b				**	***
$\Delta(1750)$	$1/2^+$	*		*		i					
$\Delta(1900)$	$1/2^-$	**	**	**		d			**	**	**
$\Delta(1905)$	$5/2^+$	****	****	*****		d			***	**	**
$\Delta(1910)$	$1/2^+$	****	**	*****		e			*	*	**
$\Delta(1920)$	$3/2^+$	***	**	***			n		***		**
$\Delta(1930)$	$5/2^-$	***		***							
$\Delta(1940)$	$3/2^-$	**	**	*	F						
$\Delta(1950)$	$7/2^+$	****	****	*****	O				***	*	***
$\Delta(2000)$	$5/2^+$	**			r						**
$\Delta(2150)$	$1/2^-$	*		*		b					
$\Delta(2200)$	$7/2^-$	*		*		i					
$\Delta(2300)$	$9/2^+$	**		**		d					
$\Delta(2350)$	$5/2^-$	*		*		d					
$\Delta(2390)$	$7/2^+$	*		*		e					
$\Delta(2400)$	$9/2^-$	**		**			n				
$\Delta(2420)$	$11/2^+$	****	*	****							
$\Delta(2750)$	$13/2^-$	**		**							
$\Delta(2950)$	$15/2^+$	**		**							

only 6  $\Delta$  resonances  
in MAID2007

# empirical parametrizations



the magnetic  $N\Delta$  form factors has a very simple form

$$G_M^*(Q^2) = 3 G_D(Q^2) e^{-0.21Q^2/\text{GeV}^2}$$

$Q^2_{\text{max}}$

10  $\text{GeV}^2$

for all other resonances we use the general form:

$$\bar{\mathcal{A}}_\alpha(Q^2) = \bar{\mathcal{A}}_\alpha(0)(1 + a_1 Q^2 + a_2 Q^4 + a_3 Q^6 + a_4 Q^8) e^{-b_1 Q^2}$$

numerical examples for a few resonances:

(complete results for protons and neutrons are found in our Review EPJ ST 198 (2011) 141)

$N^*, \Delta^*$		$\bar{\mathcal{A}}_\alpha(0)$	$a_1$	$a_2$	$a_4$	$b_1$	
$P_{11}(1440)p$	$A_{1/2}$	-61.4	0.871	-3.516	-0.158	1.36	5 $\text{GeV}^2$
	$S_{1/2}$	4.2	40.	0	1.50	1.75	
$D_{13}(1520)p$	$A_{1/2}$	-27.4	8.580	-0.252	0.357	1.20	5
	$A_{3/2}$	160.6	-0.820	0.541	-0.016	1.06	
	$S_{1/2}$	-63.5	4.19	0	0	3.40	
$S_{11}(1535)p$	$A_{1/2}$	66.4	1.608	0	0	0.70	5
	$S_{1/2}$	-2.0	23.9	0	0	0.81	
$D_{15}(1675)p$	$A_{1/2}$	15.3	0.10	0	0	2.00	??
	$A_{3/2}$	21.6	1.91	0.18	0	0.69	
	$S_{1/2}$	1.1	0	0	0	2.00	
$F_{15}(1680)p$	$A_{1/2}$	-25.1	3.780	-0.292	0.080	1.25	4
	$A_{3/2}$	134.3	1.016	0.222	0.237	2.41	
	$S_{1/2}$	-44.0	3.783	0	0	1.85	

# empirical parametrizations for large $Q^2$



the Maid parametrization with Gaussian forms for large  $Q^2$   
is convenient and leads to fewer terms

However, it violates pQCD, which predicts:  $A_{1/2}(Q^2) \sim 1/Q^3$

$$A_{3/2}(Q^2) \sim 1/Q^5$$

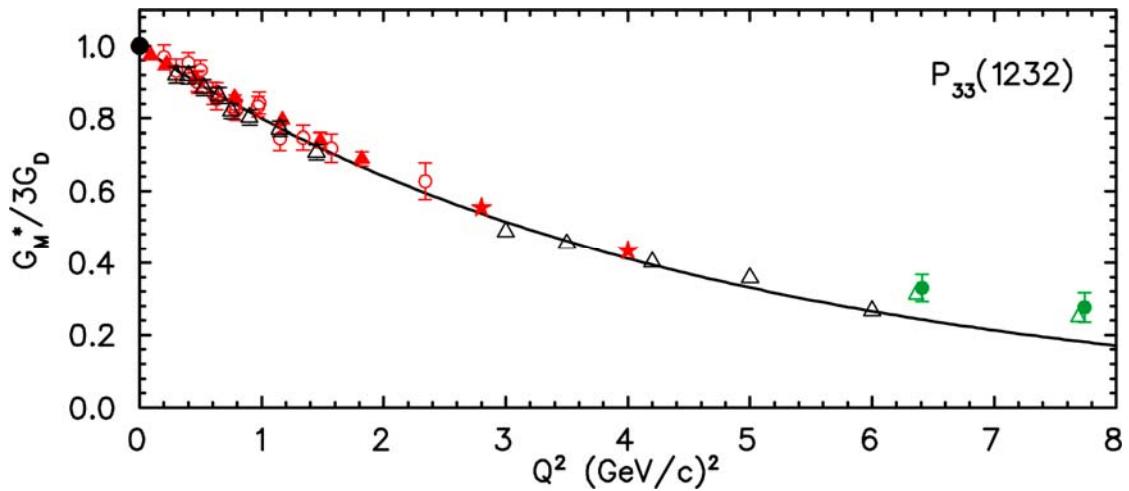
$$S_{1/2}(Q^2) \sim 1/Q^3$$

improved ansatz:

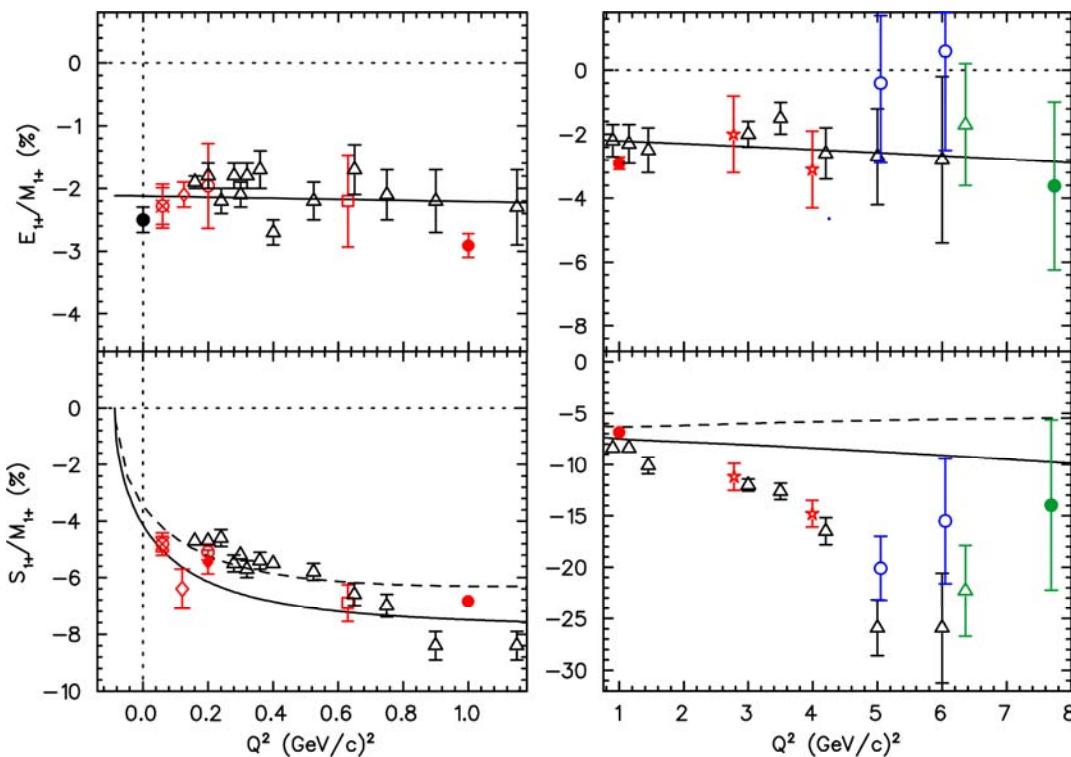
$$A_{1/2}, S_{1/2} \sim \frac{k(Q^2, M^*)}{k(0, M^*)} \frac{G_{dip}(Q^2)}{\sqrt{1 + \frac{Q^2}{(M \pm M^*)^2}}} \frac{P_n(Q^2)}{Q_n(Q^2)} \xrightarrow{Q \rightarrow \infty} \frac{1}{Q^3}$$

$$A_{3/2} \sim \frac{k(Q^2, M^*)}{k(0, M^*)} \frac{G_{dip}(Q^2)}{\sqrt{1 + \frac{Q^2}{(M \pm M^*)^2}}} \frac{P_{n-1}(Q^2)}{Q_n(Q^2)} \xrightarrow{Q \rightarrow \infty} \frac{1}{Q^5}$$

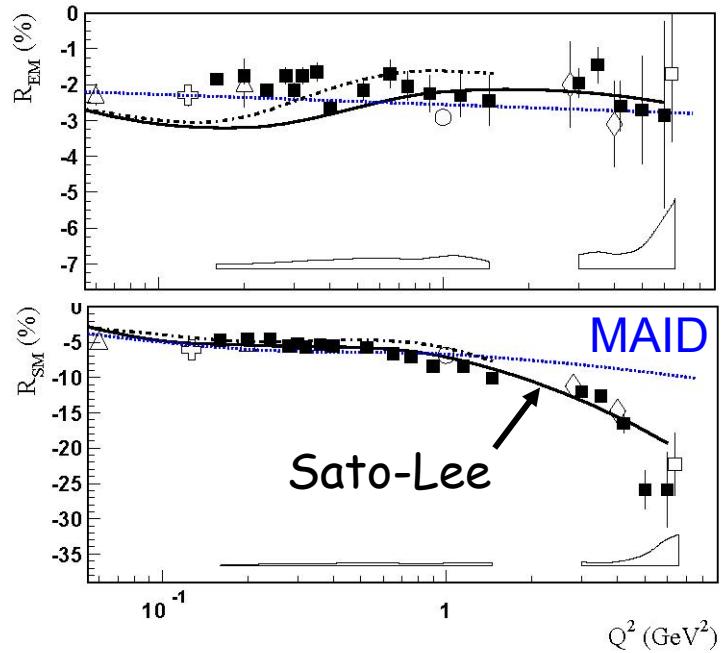
# $N \rightarrow \Delta(1232)$ transition form factors



● ● ● MAID analysis  
△ △ JLab analysis  
○ MAID analysis revisited  
 for narrow energy range  $\sim 1232$  MeV



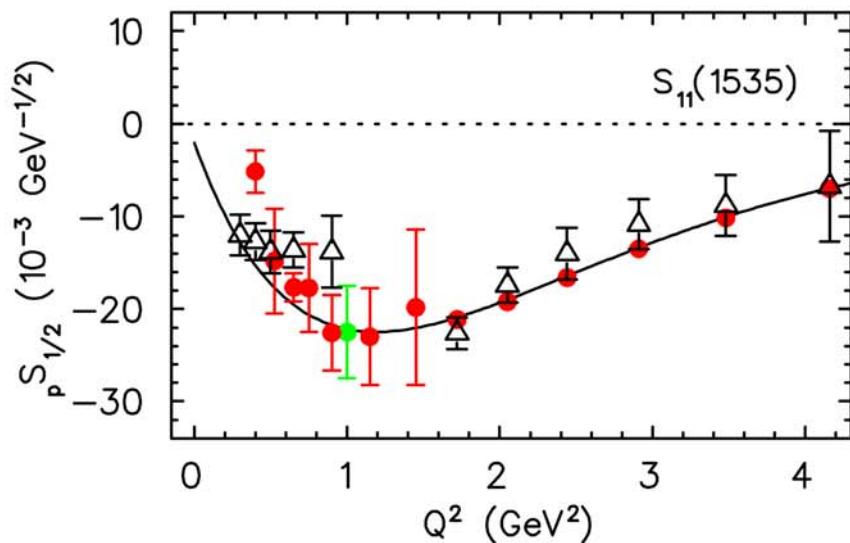
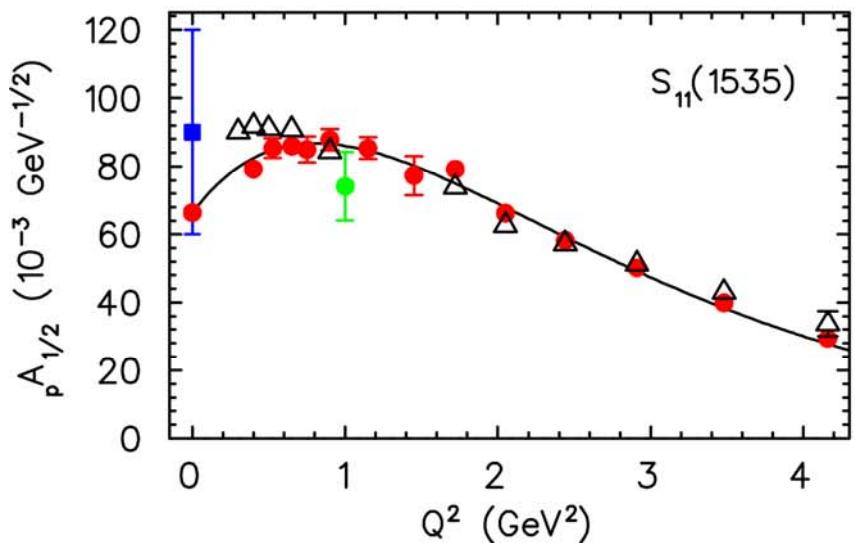
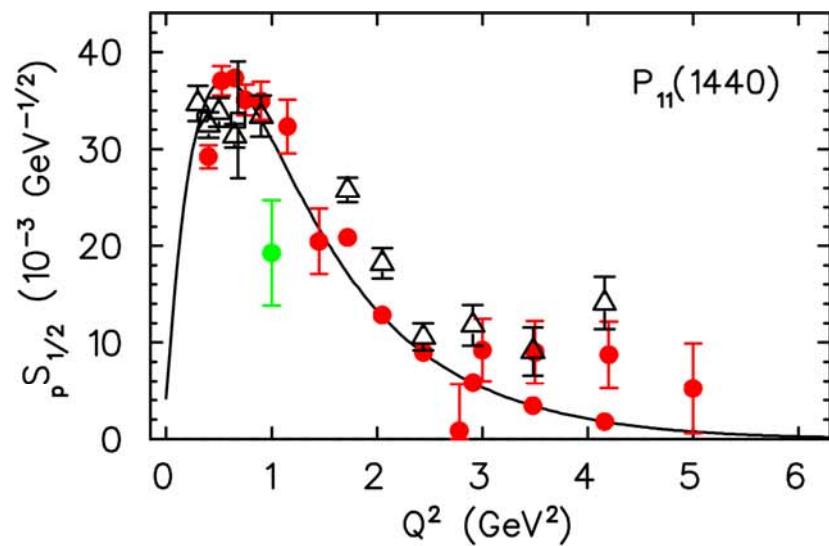
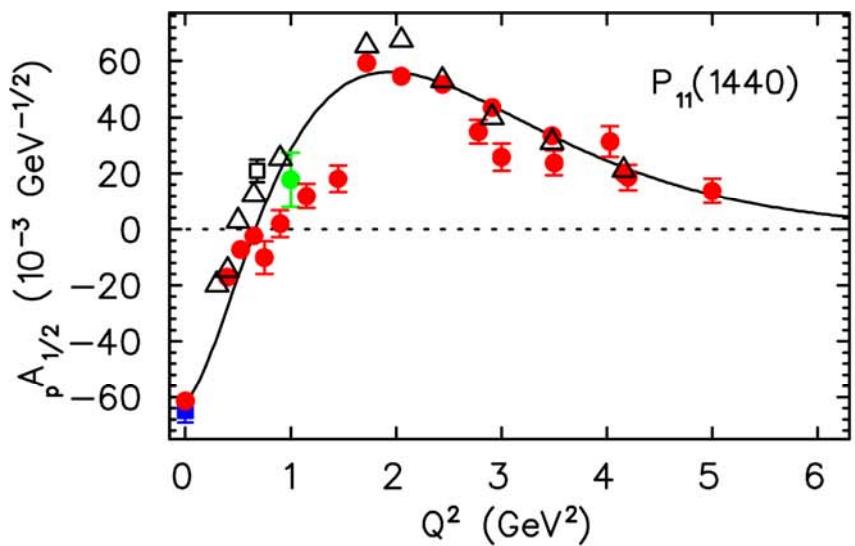
one of few cases with disagreement between Mainz and JLab analysis



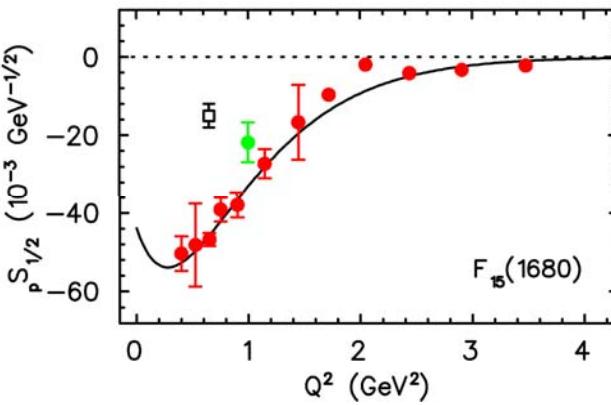
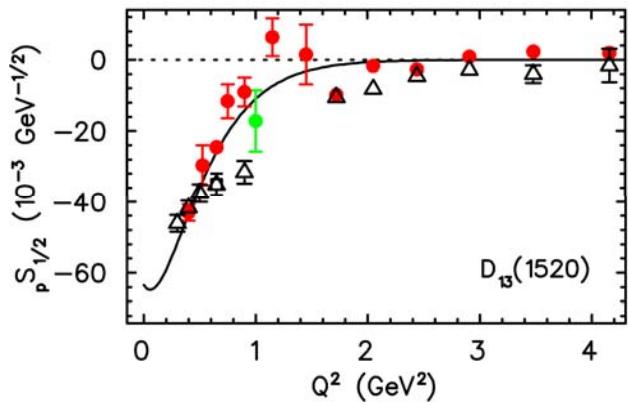
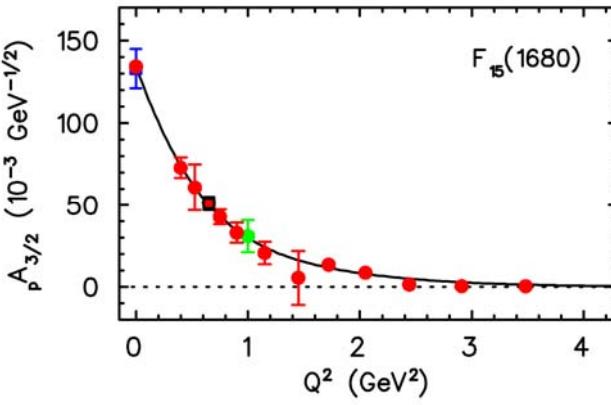
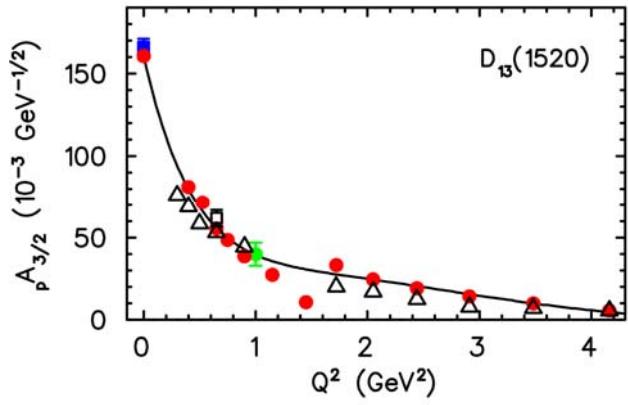
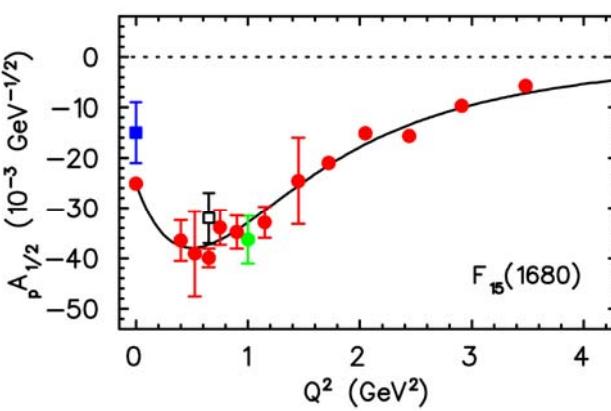
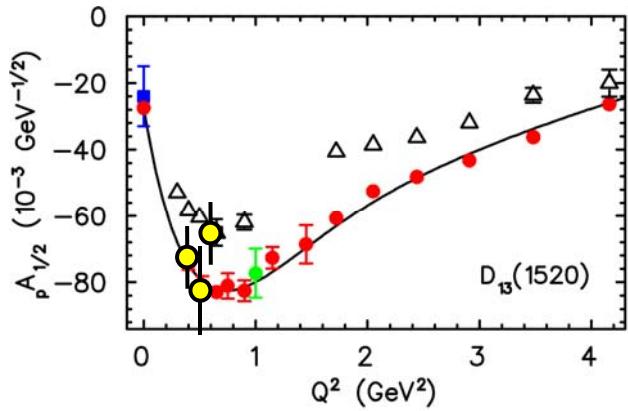
# $N \rightarrow N(1440)1/2^+$ and $N \rightarrow N(1535)1/2^-$ excitation



from MAID and JLab analyses



# transition FFs for $N \rightarrow N(1520)3/2^-$ and $N \rightarrow N(1680)5/2^+$ excitation



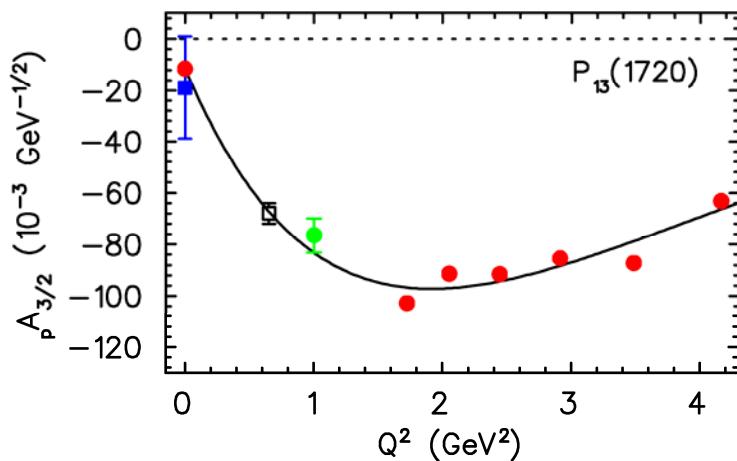
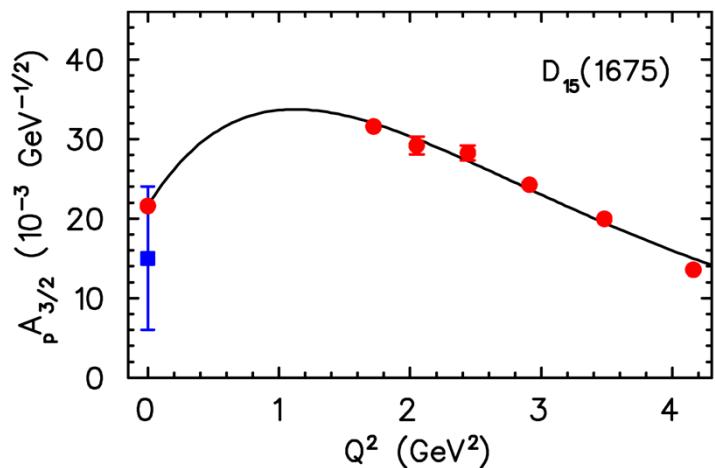
data :  
 practically all  
 underlying cross sections  
 that went into the fits  
 are from CLAS

analysis :  


---

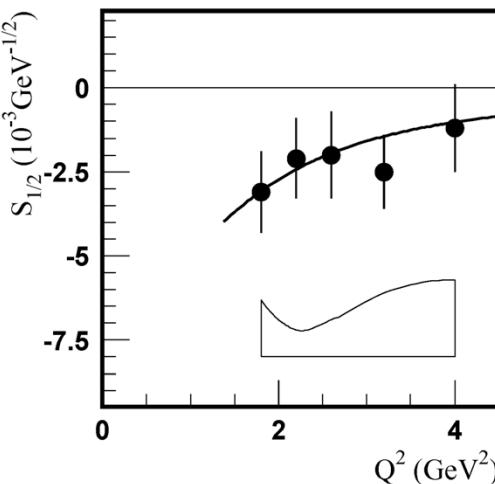
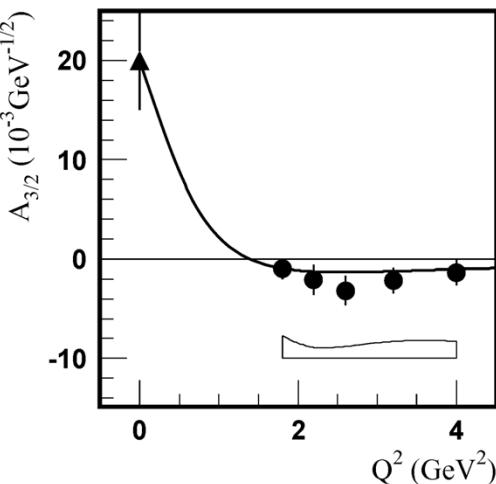
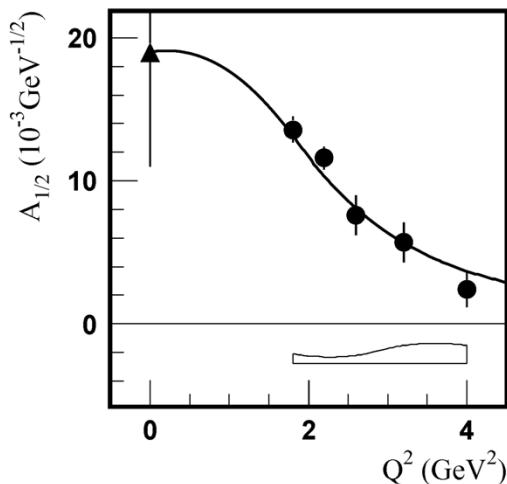
 — MAID  
 ● MAID  
 △ □ JLab  
 ○ JLab  $\pi\pi$

# $N \rightarrow N(1675)5/2^-$ and $N \rightarrow N(1720)3/2^+$ excitations



corresponding  $A_{1/2}$  and  $S_{1/2}$  ff were found to be small, consistent with zero

JLab analysis (V. Burkert et al.), new in PDG2015: here e.g.  $N(1675)5/2^-$   $D_{15}$



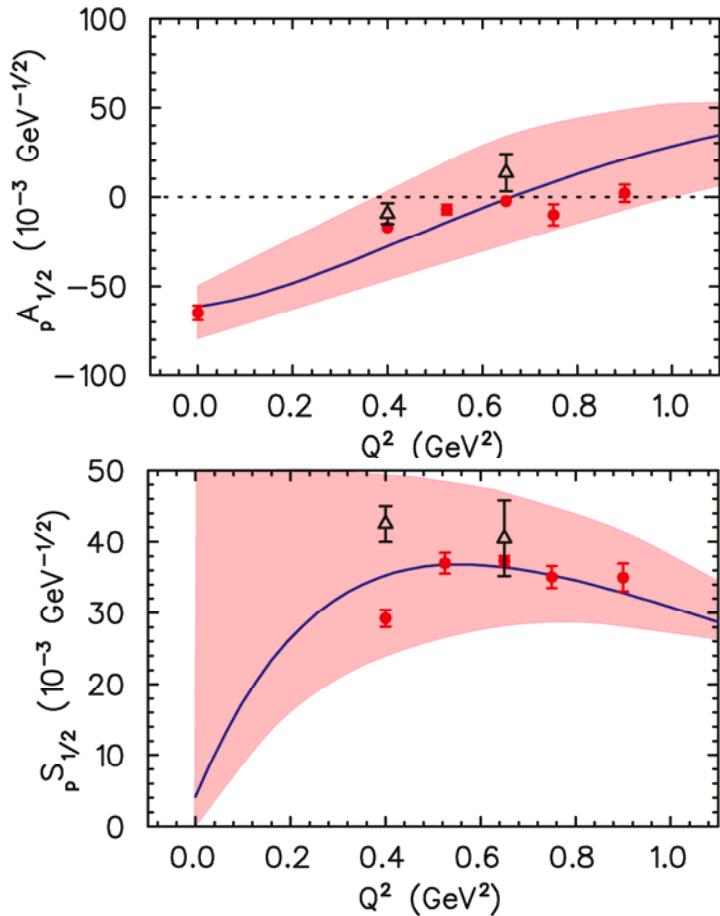
# Form factor constraints at low $Q^2$



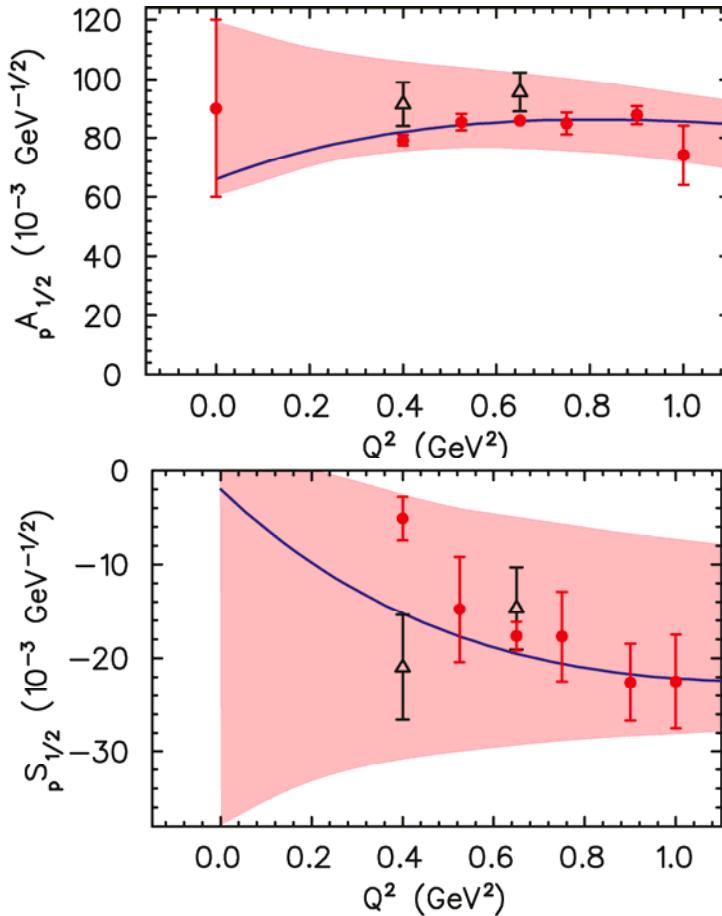
transverse form factors are constrained by photoproduction

longitudinal form factors are un-constrained ???

$N \rightarrow N(1440) 1/2^+$



$N \rightarrow N(1535) 1/2^-$



pink error bands are not to be quoted, artistic view

# Long-Wavelength Limit (Siegert Theorem, $k_R \rightarrow 0$ )



In the physical region the 3-momentum of the virtual photon at  $W=M_R$  is always finite (positive)

However, the transition form factors can be extrapolated into the unphysical region down to the pseudo-threshold, where  $k_R = 0$

The pseudo-threshold is defined by

$$Q_{pt}^2 = -(W_R - M)^2$$

for  $\Delta(1232)$ :  $-0.087 \text{ GeV}^2$ , for  $N(1440)$ :  $-0.25 \text{ GeV}^2$ , for  $N(1520)$ :  $-0.34 \text{ GeV}^2$

In the LWL longitudinal and electric multipoles are no longer independent:

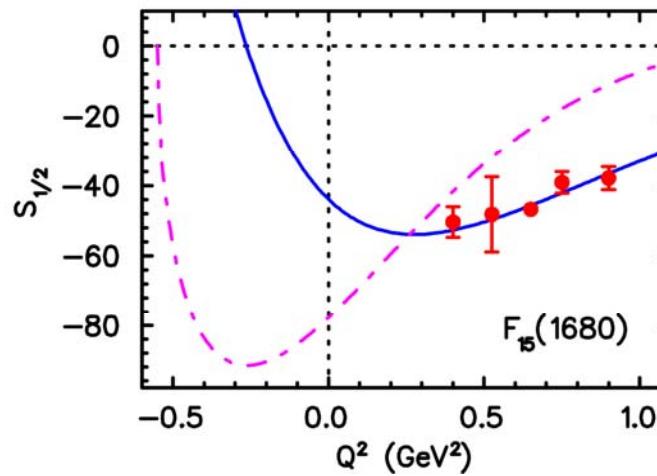
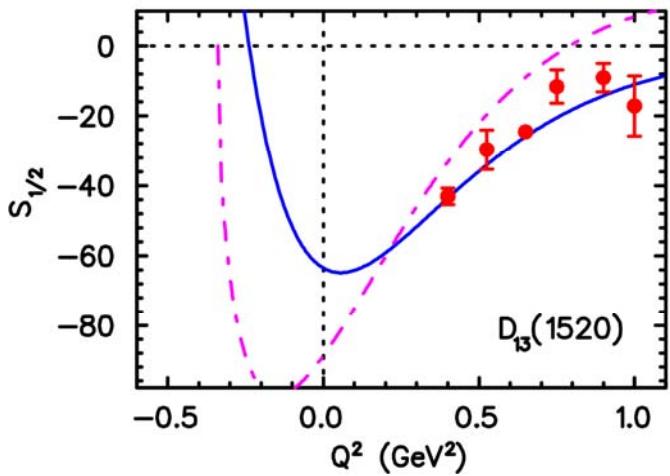
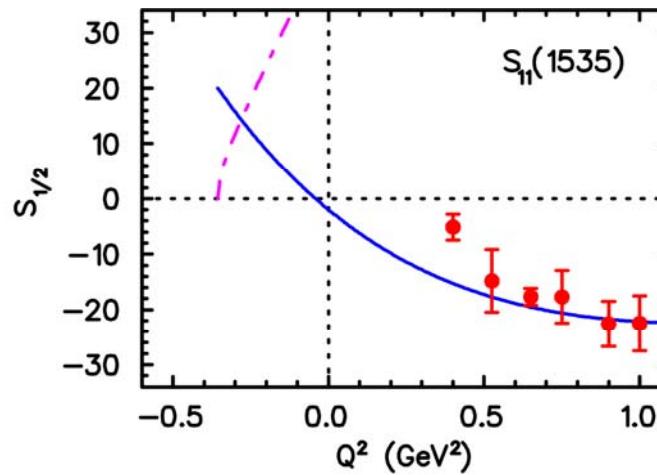
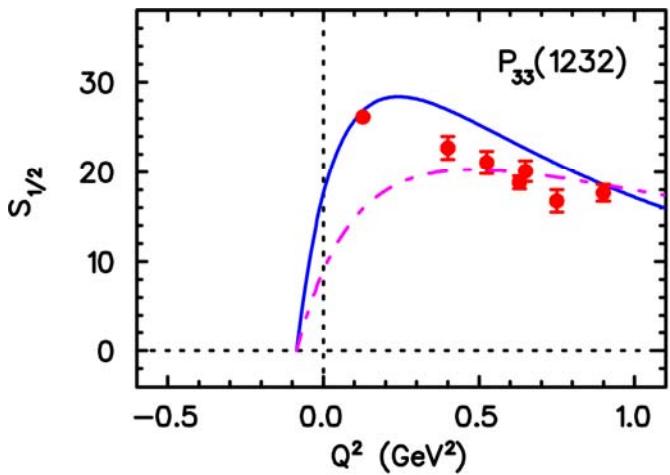
$$\begin{aligned} S_{\ell+} &= \frac{k}{\omega} L_{\ell+} \quad \rightarrow \quad \frac{k}{\omega} E_{\ell+} \quad (l \geq 0) \\ S_{\ell-} &= \frac{k}{\omega} L_{\ell-} \quad \rightarrow \quad -\frac{\ell-1}{\ell} \frac{k}{\omega} E_{\ell-} \quad (l \geq 2) \end{aligned}$$

but there is no such restriction for the Roper  $S_{1-}$



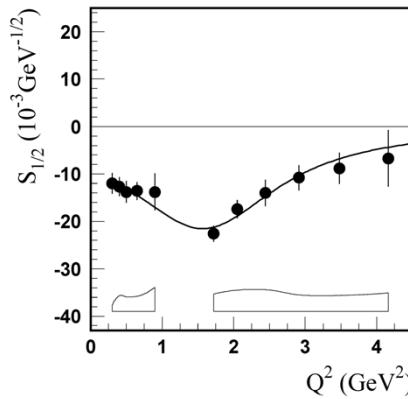
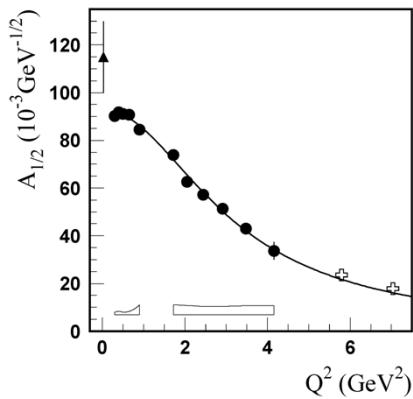
# transition form factors at low $Q^2$

— empirical fits MAID2007  
- - - Siegert Theorem (long-wavelength limit)



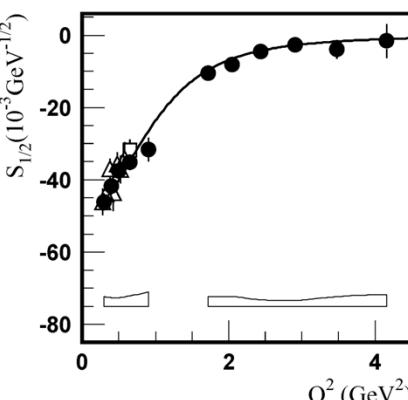
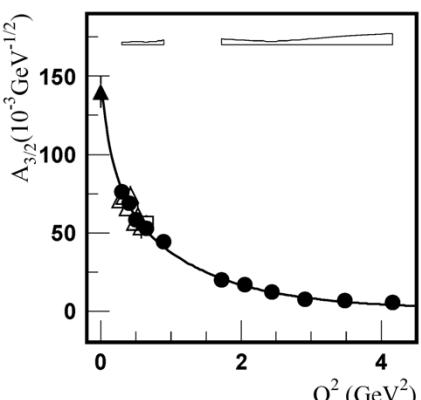
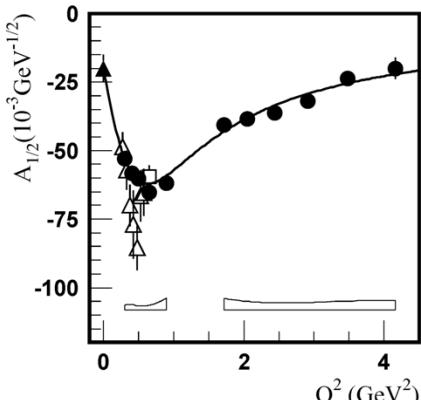
problem with  $S_{11}$   
violation of LWL ?

# JLab analysis of transition form factors in PDG 2014



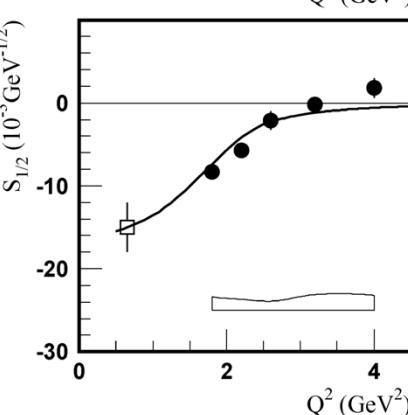
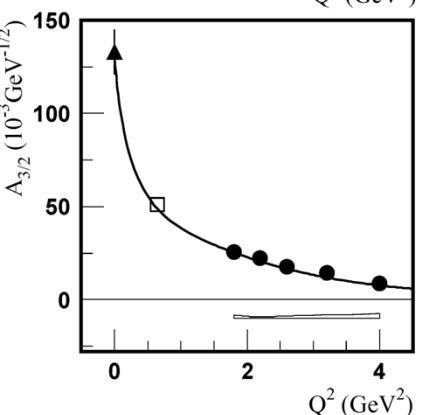
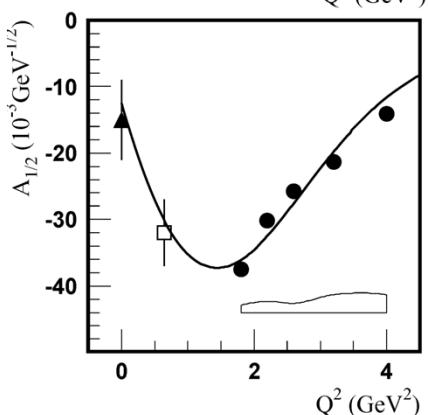
N(1535)1/2<sup>-</sup> S11

$$S_{1/2}(Q^2) < 0$$



N(1520)3/2<sup>-</sup> D13

$$S_{1/2}(Q^2) < 0$$



N(1680)5/2<sup>+</sup> F15

$$S_{1/2}(Q^2) < 0$$

# Outlook



## work in progress

V. Kashevarov, L.T.

EtaMAID ( $\gamma, \eta$  and  $\gamma, \eta'$ ) with Regge and without  
and search for  $\eta$  and  $\eta'$  couplings in all N resonances

Mainz-Tuzla-Zagreb Collaboration

model-independent single-energy PWA with fixed-t analyticity for  $\gamma, \eta$

A. Svarc, R. Workman, L.T.

transition form factors at the pole for MAID and SAID solutions  
with L+P expansion method in  $Q^2$  region:  $0 < Q^2 < 5 \text{ GeV}^2$

## work planned for 2016

V. Kashevarov, L.T.

MAID ( $e, e'\pi$ ) update with 2- and 3-star N and  $\Delta$  resonances