



Exclusive Single Pion off the Proton: Results from CLAS

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The most challenging problems in Hadron Physics

- Non-perturbative **DCSB** generates more than 98% of dress quark masses as well as dynamical structure
 - although, higgs mechanism $< 2\%$ in N , N^* masses
- Quark-gluon confinement in baryons emerges from QCD
 - dressed quarks, meson-baryon cloud, dressed gluon,...

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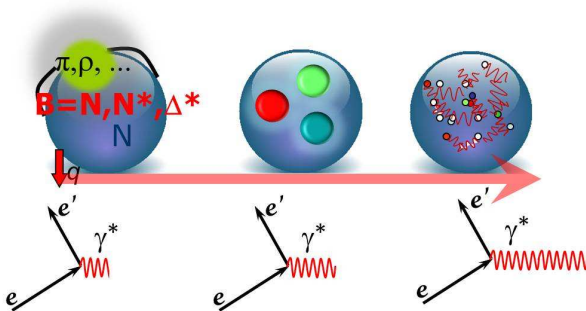
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The most challenging problems in Hadron Physics

- Study of the excited states of the nucleon is important step in the development of a fundamental understanding of strong interaction

[N. Isgur, V. Burkert (2000)]

- The most fundamental question: "WHAT ARE THE RELEVANT DEGREE-OF-FREEDOM AT VARYING DISTANCE SCALE ?"

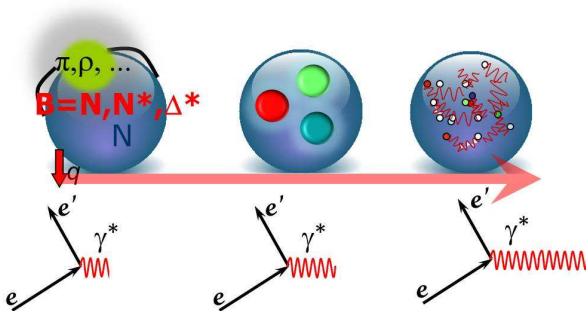


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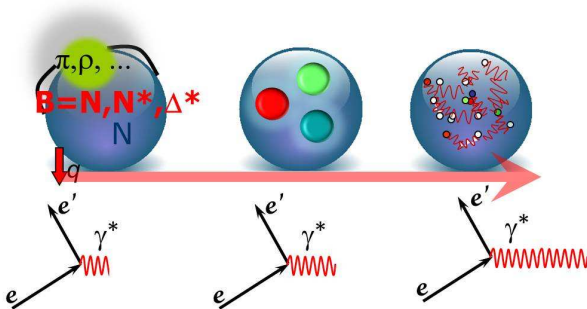


The most challenging problems in Hadron Physics

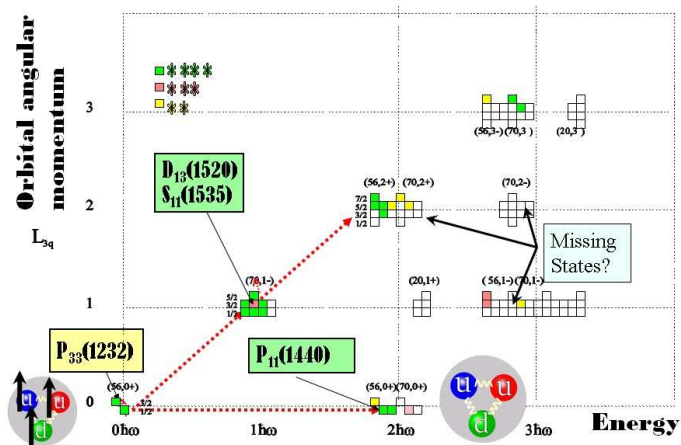
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- The most fundamental question: “ **WHAT ARE THE RELEVANT DEGREE-OF-FREEDOM AT VARYING DISTANCE SCALE ?** ”



$SU(6) \times O(3)$ Classification of Baryons

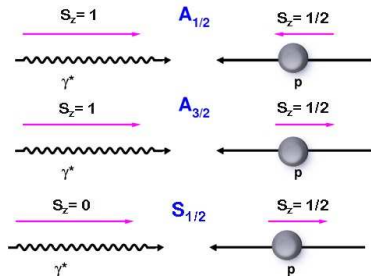


- There are questions about underlying DoF of some well known state...but still many open questions.. related with QCD, FT, CQM, LQCD ...

Photo-coupling Amplitudes

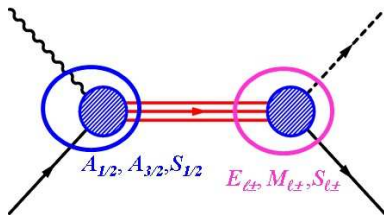
Spin combination

- Transverse
- Longitudinal

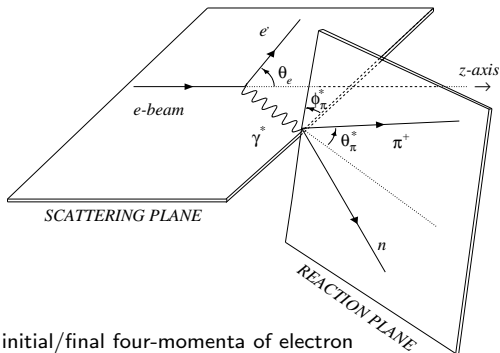


Multipole Amplitudes

- $E_{l\pm}$, $M_{l\pm}$, and $S_{l\pm}$
 - l : the orbital angular momentum in $N\pi$ system
 - \pm sign: spin of proton couples to the orbital momentum

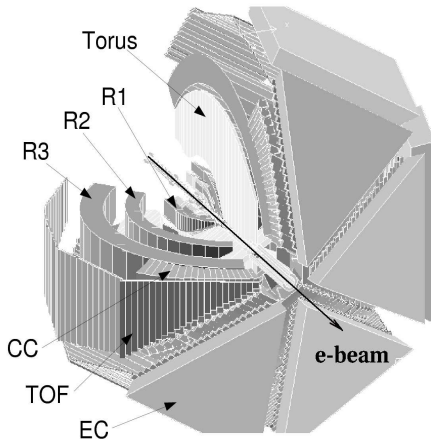
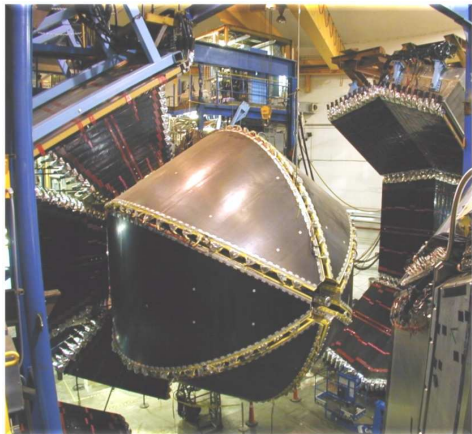


Reaction



- $k_{i,f}$: the initial/final four-momenta of electron
- $E_{i,f}$: the initial/final energy of electron
- θ_e : the electron scattering angle
- $p_{\gamma,i}$: the virtual photon/target four-momenta
- $W^2 = (p_\gamma + p_i)^2 = M_p^2 + 2M_p\nu - Q^2$
- ν : transferred energy $= E_i - E_f = \frac{p_i \cdot p_\gamma}{M_p}$
- Q^2 : virtuality of the exchanged photon $= -(k_i - k_f)^2 = 4E_i E_f \sin^2(\theta_e/2)$
- θ_π^* : the angle between the virtual photon and the hadron (π^+)
- ϕ_π^* : the angle between the electron scattering plane and the hadronic production plane

CEBAF Large Acceptance Spectrometer



CLAS single pion data coverage

Final State	W (GeV)	Q^2 (GeV ²)	Observables
$n\pi^+$	1.1 -1.38	0.16-0.36	$d\sigma/d\Omega$
	1.1 -1.55	0.3 -0.6	$d\sigma/d\Omega$
	1.1 -1.7	1.7 -4.5	$d\sigma/d\Omega, A_B$
	1.65-2.0	1.8 -4.5	$d\sigma/d\Omega$
$p\pi^0$	1.1 -1.38	0.16-0.36	$d\sigma/d\Omega$
	1.1 -1.68	0.4 -1.8	$d\sigma/d\Omega, A_B, A_T, A_{BT}$
	1.1 -1.39	3.0 -6.0	$d\sigma/d\Omega$

- All CLAS data is available at CLAS-DB
<http://clasweb.jlab.org/physicsdb/>

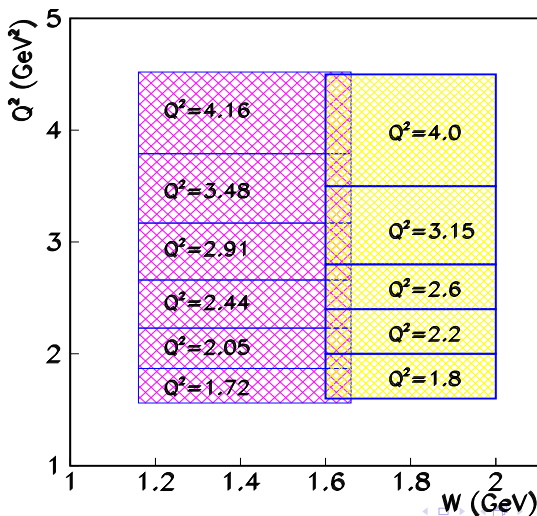
[** $K\Lambda, K\Sigma$ see talk by D. Carman, 2π see talk by V. Mokeev]

CLAS single pion data coverage for $n\pi^+$

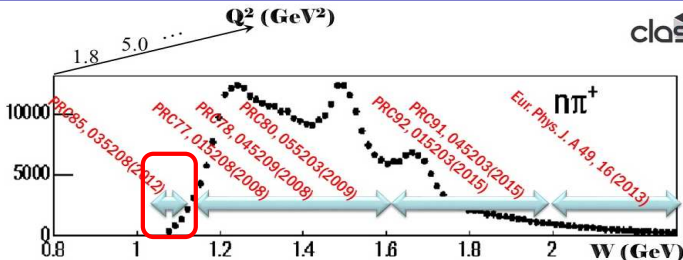
RunGroup	W (GeV)	Q^2 (GeV ²)	Observables	# data
e1-6	1.10 -1.15	1.8 -4.5	$d\sigma/d\Omega$	1800
e1-6	1.1 -1.7	1.7 -4.5	$d\sigma/d\Omega, (A_{LU})$	50400 (12600)
e1-f	1.65 -2.0	1.8 -4.5	$d\sigma/d\Omega$	32500
e1-6	0.16-0.58 (x_{BJ})	1.6 -4.5	$d\sigma/dt$	140

Kinematic coverage of data

- Kinematic range W , Q^2 of N^* analyses
- ($\Delta(1232)3/2^+$), $N(1440)1/2^+$, $N(1520)3/2^-$, $N(1535)1/2^-$, $N(1675)5/2^-$, $N(1680)5/2^+$, and $N(1710)1/2^+$

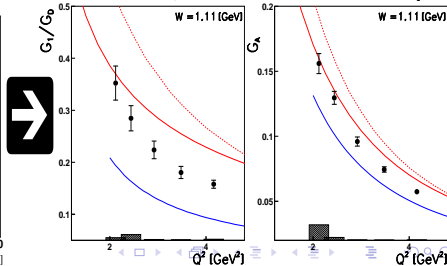
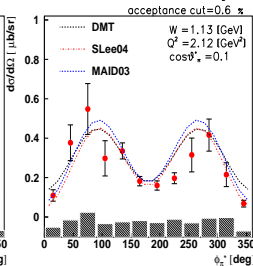
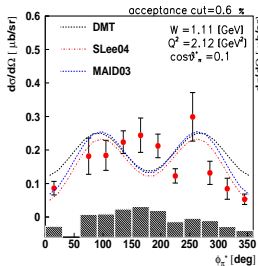


$\vec{e}p \rightarrow e'\pi^+n$ near pion threshold ($1.1 < W < 1.15$ GeV)

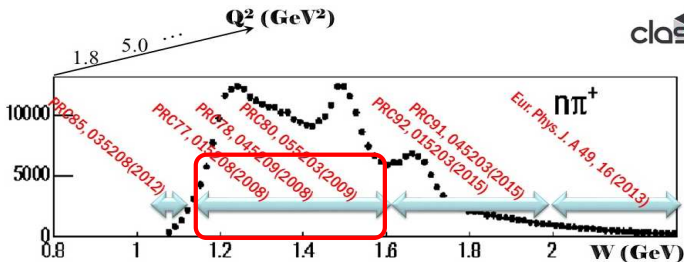


- Generalized form factor (G_1) and Axial Form Factor (G_A) near pion threshold
- Multipole fit vs. LCSR, Both are consistent result in lowest W
- Transverse s-wave multipole (E_{0+}) is dominated

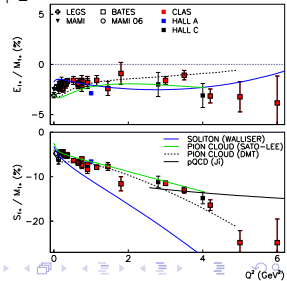
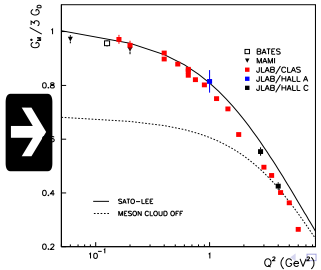
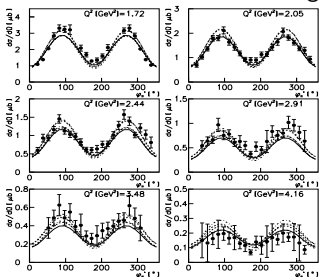
[red solid: LCSR+FF, dash: pure LCSR, blue solid: MAID07 ↓]



$\vec{e}p \rightarrow e'\pi^+n$ for low lying Δ ($1.15 < W < 1.69$ GeV)



- Transition Form Factors for $\Delta(1232)$... sensitive with π^0 data
- Significant meson cloud effect in G_M^* at lower Q^2
- No indication of leading pQCD contribution in $R_{EM} \rightarrow +1$



- Two different approaches: **UIM, DR**

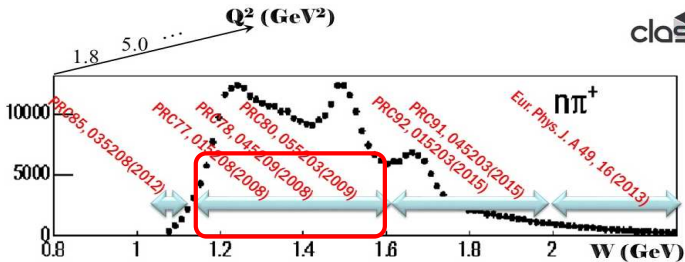
UIM

- BG UIM is built from nucleon exchange in s -, u - and π, ω, ρ exchange in t - channel
- Unitarization of multipole amplitudes in the K -matrix approximation
- Resonance contributions are parameterized in the unified BW form with energy dependence

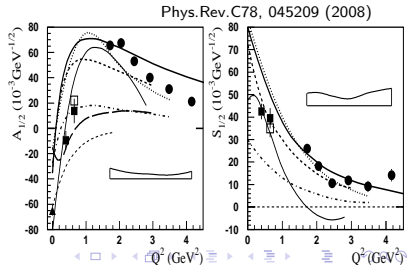
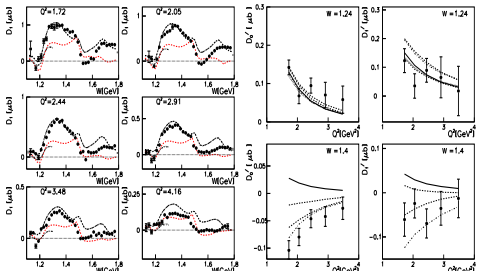
DR

- Fixed- t dispersion relation for the invariant amplitude
- Re -Amplitude to Born-term (s, u , channel nucleon exchange, π exchange in t -
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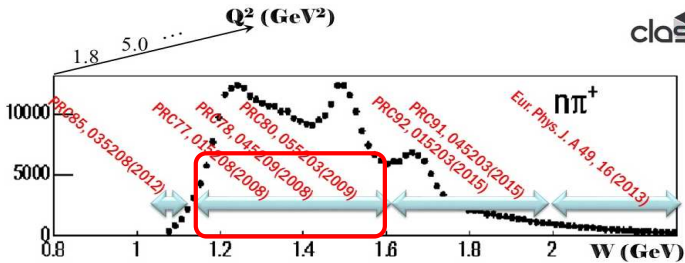
$\vec{e}p \rightarrow e'\pi^+n$ for low lying N^* ($1.15 < W < 1.69$ GeV)



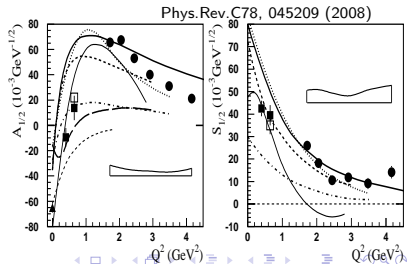
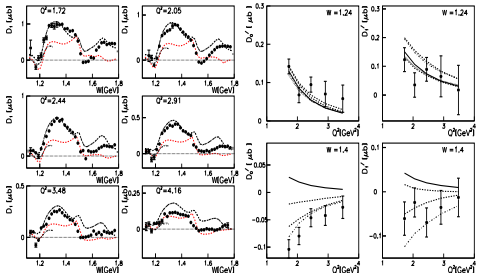
- Transition Form Factors for $N(1440)1/2^+$ (old conv: $P_{11}(1440)$)
- $A_{1/2}$ shows a sign change in $Q^2 \sim 0.8$ GeV²
- $S_{1/2}$ is large at low Q^2 and drop off smoothly with increasing Q^2



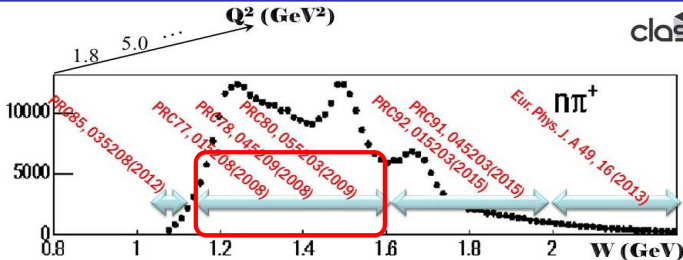
$\vec{e}p \rightarrow e'\pi^+n$ for low lying N^* ($1.15 < W < 1.69$ GeV)



- NR-Quark VM(thin-solid)/Rule out a hybrid baryon model(gluon Excite)
- LFRQ model needs a *MB* interaction at large distance
- A complex interplay btw inner core of quarks in the first radial excitation and external *MB* cloud

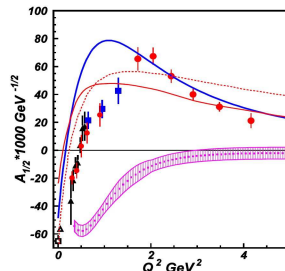
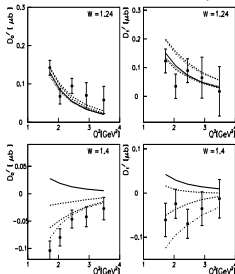
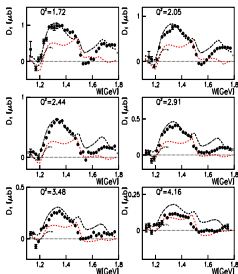


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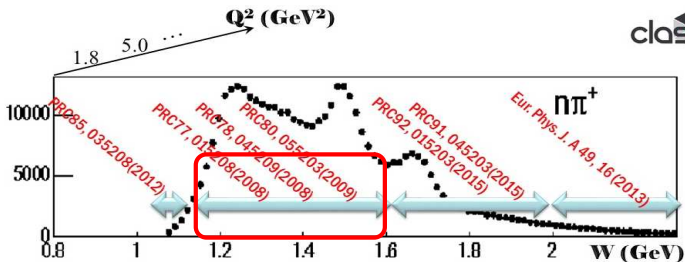


- Quark core in DSEQCD (thick blue curve), MB cloud contribution (purple band)
- $N\pi$ loops MB, running quark mass (red solid curve)
- $N\sigma$ loops MB, fixed constituent quark mass (red dashed curve)

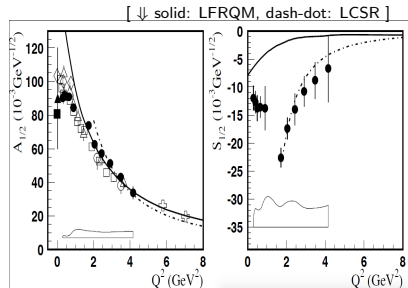
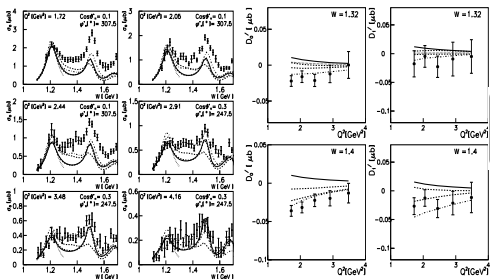
[↓ including single π and 2π data]



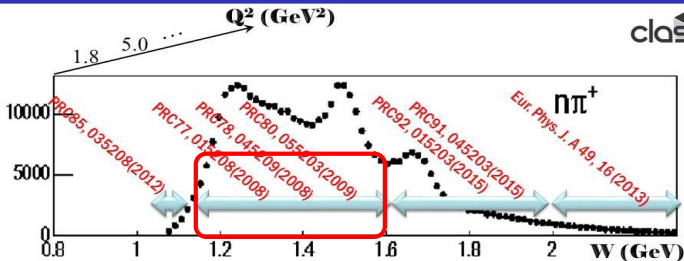
$\bar{e}p \rightarrow e'\pi^+n$ for low lying N^* ($1.15 < W < 1.69$ GeV)



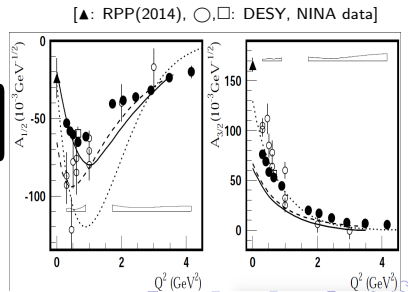
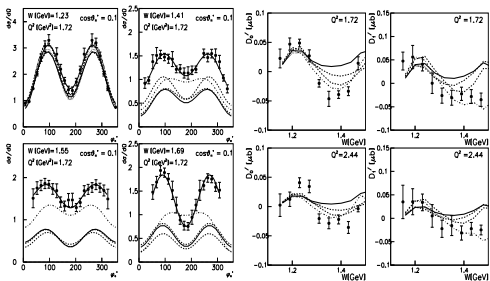
- Transition Form Factors for $N(1535)1/2^-$ (old conv: $S_{11}(1535)$)
- $\beta_{N\eta}^{PDG} = 0.45 - 0.60 \rightarrow \beta_{N\pi}^{PDG} = 0.485$ & $\beta_{N\eta}^{PDG} = 0.460$, excellent agreement $N\pi, N\eta$
- Sensitive to long. as well (strong interference $S_{11}-P_{11}$)



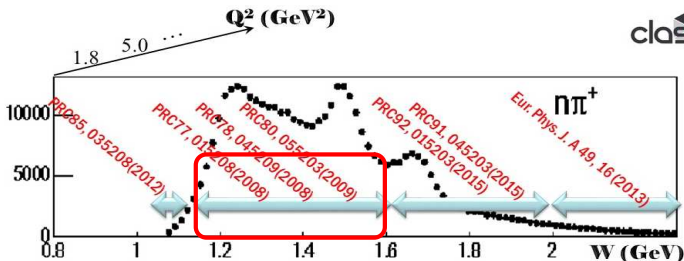
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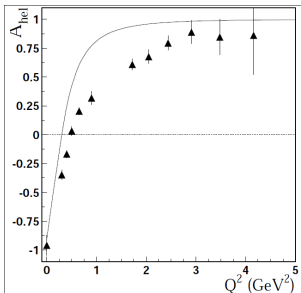
- Transition Form Factors for $N(1520)3/2^-$ (old conv: $D_{13}(1520)$)
- $A_{1/2}$ is large at high Q^2 , $A_{3/2}$ is small at high Q^2



$\vec{e}p \rightarrow e'\pi^+n$ for low lying N^* ($1.15 < W < 1.69$ GeV)



- Helicity Asymmetry for $N(1520)3/2^-$ (old conv: $D_{13}(1520)$)

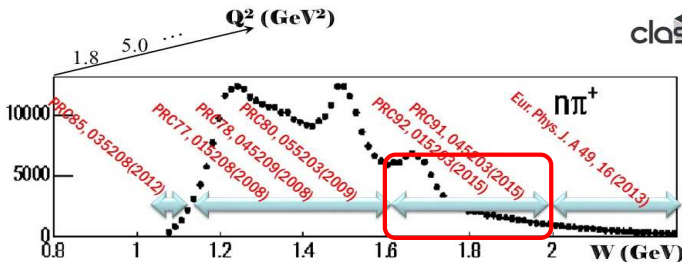


$$\frac{A_{1/2}^{D_{13}}}{A_{3/2}^{D_{13}}} = \frac{-1}{\sqrt{3}} \left(\frac{Q^2}{\alpha} - 1 \right)$$

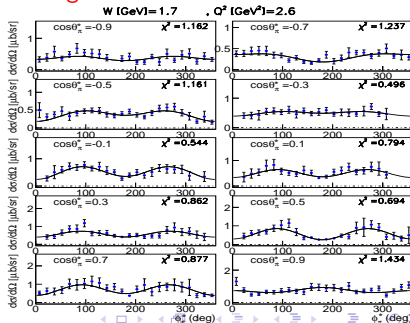
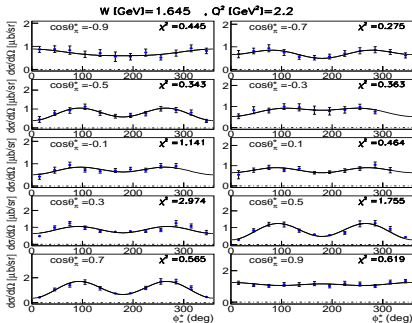
$$A_{hel} = \frac{A_{1/2}^2 - A_{3/2}^2}{A_{1/2}^2 + A_{3/2}^2}$$

- Asymptotic Q^2 behavior of A_{hel} vs. Q^2
- NRQ simple harmonic oscillator model (solid line) with spin, orbit flip amplitudes
- $A_{1/2} \ll A_{3/2}$ at low Q^2 , $A_{3/2} \ll A_{1/2}$ at high Q^2

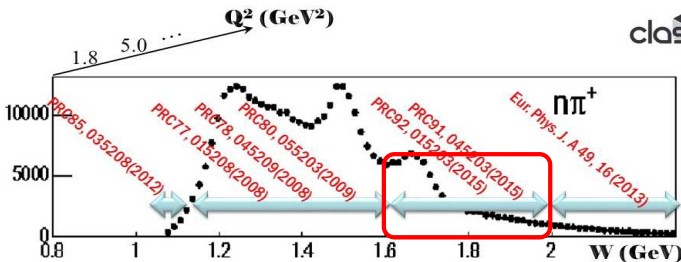
$\bar{e}p \rightarrow e'\pi^+n$ for high lying N^* ($1.65 < W < 2.0$ GeV)



• differential cross-sections for third resonance region

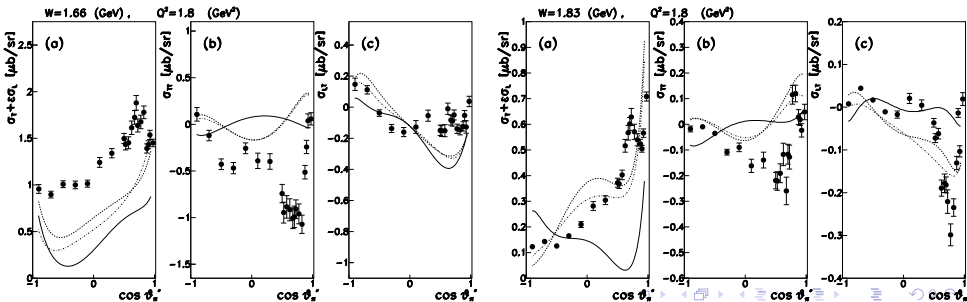


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● **structure functions ((a) $\sigma_T + \epsilon\sigma_L$, (b) σ_{TT} , (c) σ_{LT}) for third resonance region**

- dash: MAID07, dash-dot: MAID03, dot: DMT



- Two different approaches: **UIM, DR**

UIM

- BG UIM is built from nucleon exchange in s -, u - and π, ω, ρ exchange in t - channel
- Unitarization of multipole amplitudes in the K -matrix approximation
- Resonance contributions are parameterized in the unified BW form with energy dependence

DR

- Fixed- t dispersion relation for the invariant amplitude
- Re -Amplitude to Born-term (s, u , channel nucleon exchange, π exchange in t -
- Integral Im -Amplitude with the isospin structure

- **Two model-uncertainties**

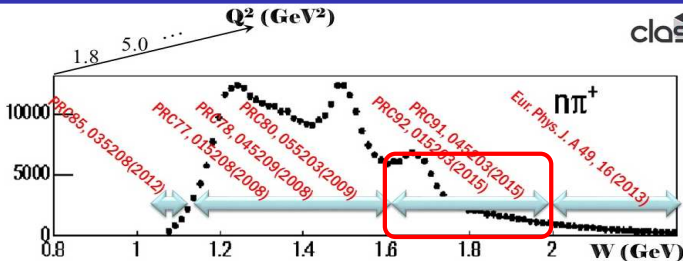
- 1/ BG determination in the UIM and Born term in DR
- 2/ A width and mass of resonances from PDG

- **Take into account...**

- 1/ All(13) **** and *** states in the 1st, 2nd, 3rd
- 2/ $\Delta(1905)F_{35}$, $\Delta(1950)F_{37}$ in 4th resonance region

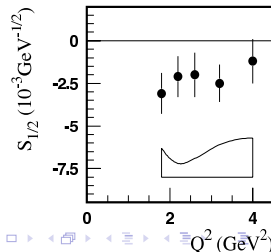
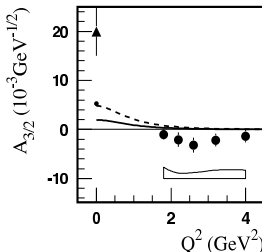
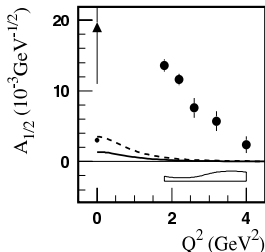
- Same BR from PDG2012

$\vec{e}p \rightarrow e'\pi^+n$ for high lying N^* ($1.65 < W < 2.0$ GeV)

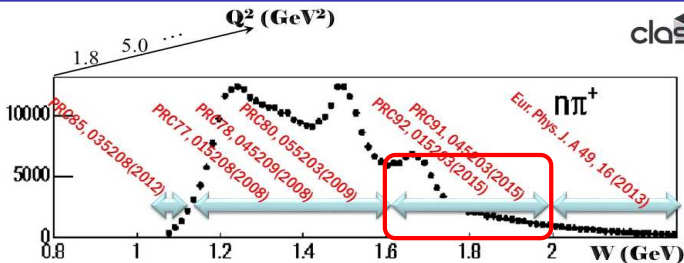


- Transition Form Factors for $N(1675)5/2^-$ (old conv: $D_{15}(1675)$)
- SQTM, Moorhouse selection rule: suppression Transverse Amplitudes

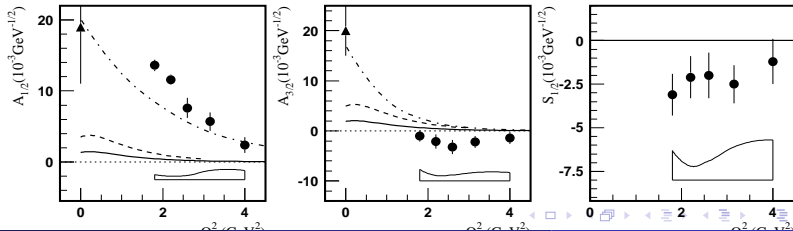
● Solid: M.M.Gianini/E.Santopinto (hQCM)
 dash: D.Merten & U.Loring(2003), Solid-dot($Q^2 = 0$): I.G.Aznauryan(LFRQ)



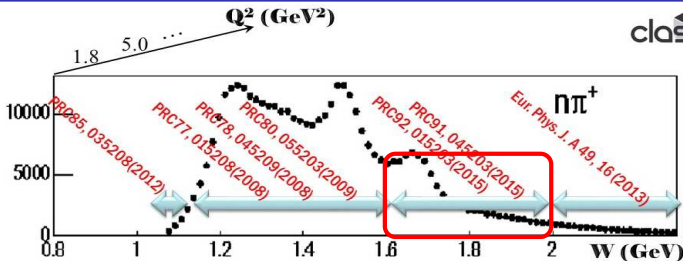
$\vec{e}p \rightarrow e'\pi^+n$ for high lying N^* ($1.65 < W < 2.0$ GeV)



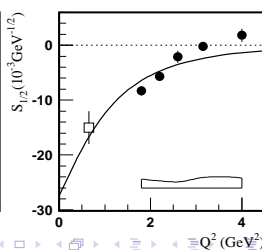
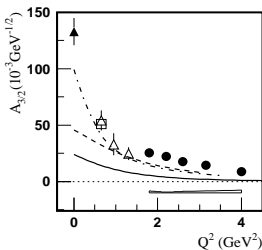
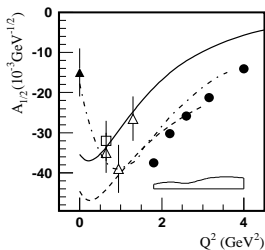
- Non-quark contributions dominance, A strong coupling $A_{1/2}$ for $Q^2 < 4$ GeV²
- Significant MB contribution from the dynamical coupled-channel model (dash-dot: B. Julia-Diaz, T-S.H. Lee, A. Matsuyama)
- A strong suppression of $A_{3/2}$ for $Q^2 > 1.8$ GeV²



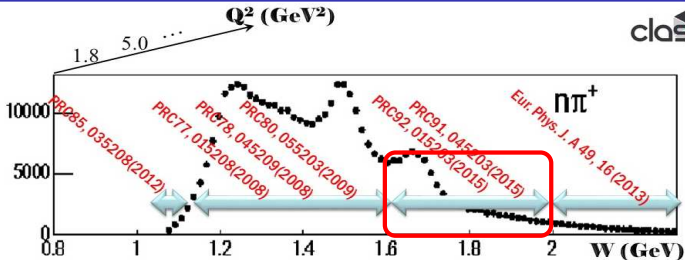
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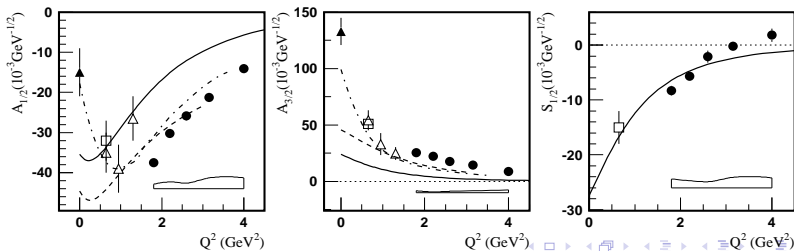
- Transition Form Factors for $N(1680)5/2^+$ (old conv: $F_{15}(1680)$)
- ▲ RPP(PDG:2014), △ V.Mokeev& I.G.Aznauryan(2013), □ I.G.Aznauryan(2005)
- Solid: M.M.Gianini/E.Santopinto (hQCM), dash-dot: Z.Lee& F.Close(1990), dash: D.Merten& U.Loring(2003)



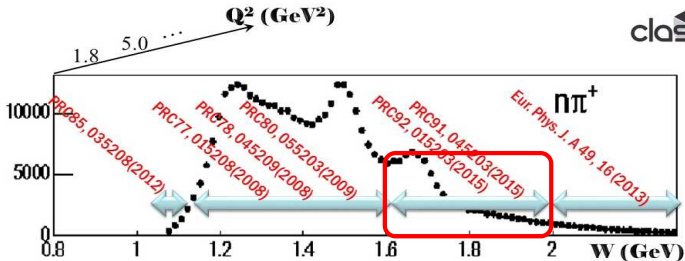
$\vec{e}p \rightarrow e'\pi^+n$ for high lying N^* ($1.65 < W < 2.0$ GeV)



- Transition Form Factors for $N(1680)5/2^+$ (old conv: $F_{15}(1680)$)
- All models estimates amplitudes larger $A_{1/2}$ (lower $A_{3/2}$) than data
- MB contribution should be taken into account ?



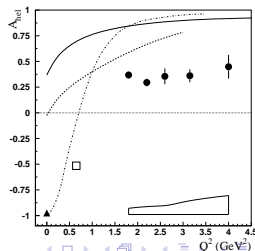
$\vec{e}p \rightarrow e'\pi^+n$ for high lying N^* ($1.65 < W < 2.0$ GeV)



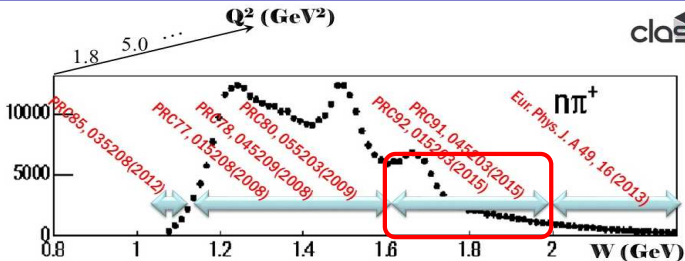
- Helicity asymmetry shows a very slow rise at $Q^2 > 2\text{GeV}^2$
- Interesting of helicity asymmetry $Q^2 > 5 \text{ GeV}^2$?
→ CLAS12

$$A_{hel} = \frac{A_{1/2}^2 - A_{3/2}^2}{A_{1/2}^2 + A_{3/2}^2}$$

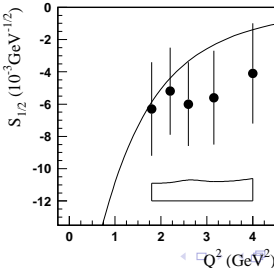
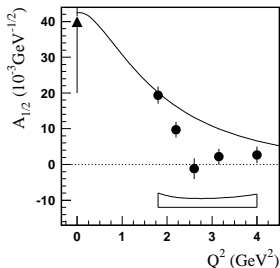
- □ CLAS single- π and 2π electroproduction
- ▲ RPP2014 at $Q^2 = 0$
- Solid: M.M.Gianini/E.Santopinto (hQCM), dash-dot: Z.Lee& F.Close(1990), dash: D.Merten& U.Loring(2003)



$\vec{e}p \rightarrow e'\pi^+n$ for high lying N^* ($1.65 < W < 2.0$ GeV)

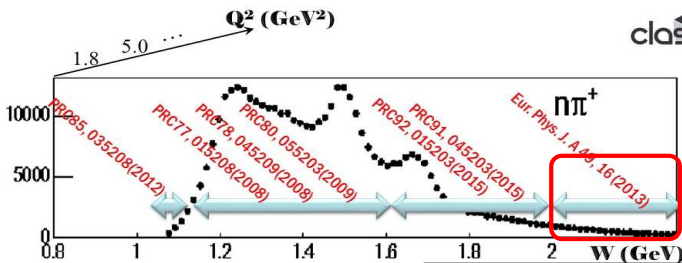


- Transition Form Factors for $N(1710)1/2^+$ (old conv: $P_{11}(1710)$)
- Finite size of $A_{1/2}$ for $Q^2 < 2.5$ GeV^2
- Finite size and negative of $S_{1/2}$ for all given Q^2 GeV^2



Solid: M.M.Gianini
E.Santopinto (hQCM)

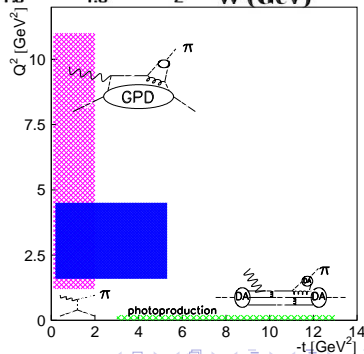
$\vec{e}p \rightarrow e'\pi^+n$ for Deep Process ($W > 2.0$ GeV)



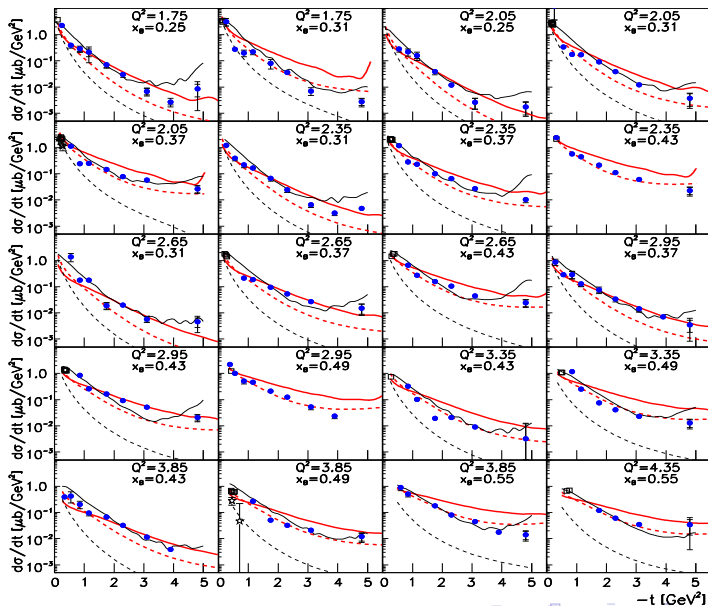
- Transition between hadronic and partonic picture of strong interaction
- Connection to the transversity GPD
- Kinematic variables $\rightarrow x_{BJ}, Q^2$ and $-t$

Blue box \rightarrow

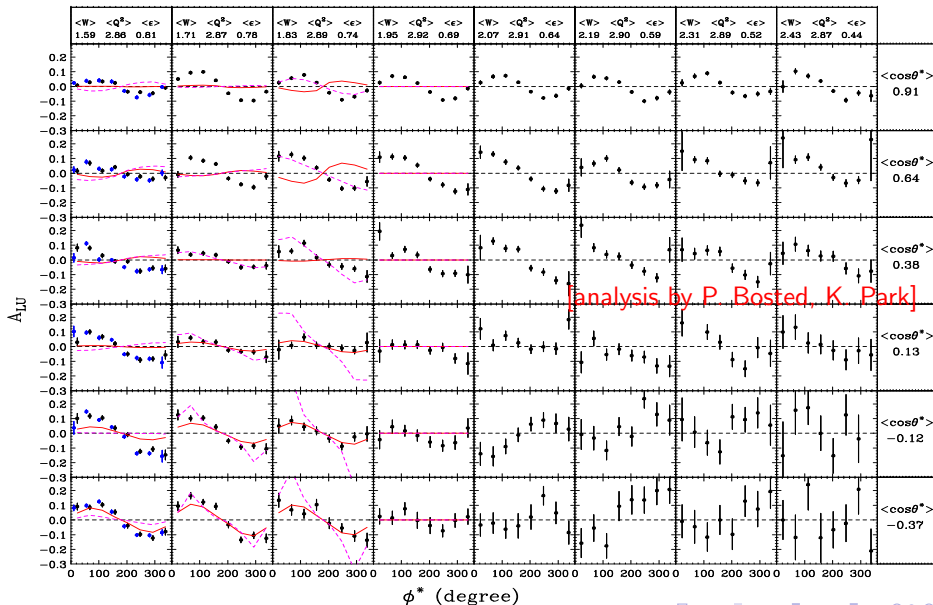
Note that $-t$: the momentum transfer to the nucleon, $t = (p_N - p'_N)^2$



$\vec{e}p \rightarrow e'\pi^+n$ for Deep Process ($W > 2.0$ GeV)



More single pion data are coming from CLAS6!



$\vec{e}p \rightarrow e'\pi^0 p$ for (low Q^2)

$E_0 = 2$ GeV, $W = 1.1-1.8$ GeV, $Q^2 = 0.4-1.0$ GeV²

[analysis by N. Markov]

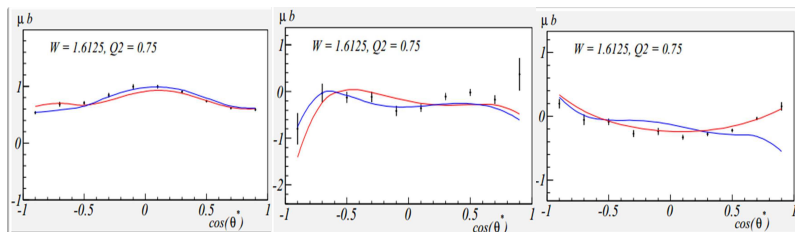


Figure: (PRELIMINARY) $\sigma_T + \epsilon\sigma_L$, σ_{TT} , and σ_{LT} , red(MAID07), blue(SAID08) curves

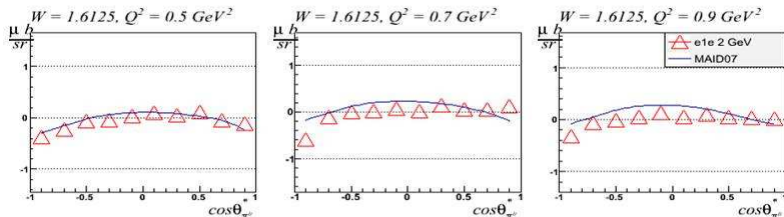


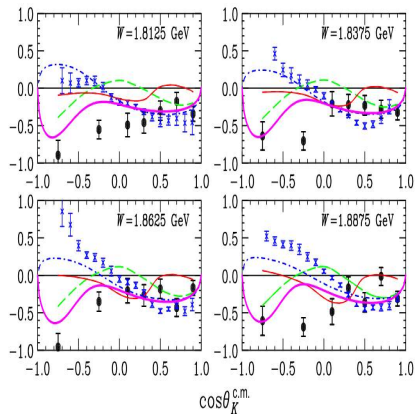
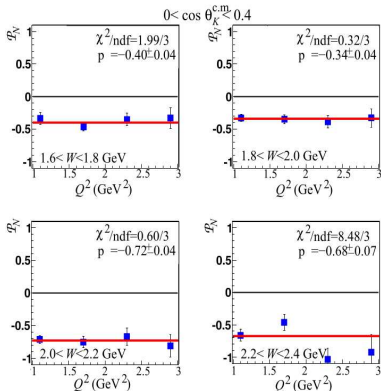
Figure: (PRELIMINARY) σ_{LT} , blue(MAID07) curve

- **CLAS6** → **CLAS12 N^* Physics Program**

- **E12-09-003**, Nucleon Resonance Studies with CLAS12
- **E12-06-108A**, *KY* electroproduction with CLAS12
- **LOI12-15-004**, Search for Hybrid Baryons with CLAS12

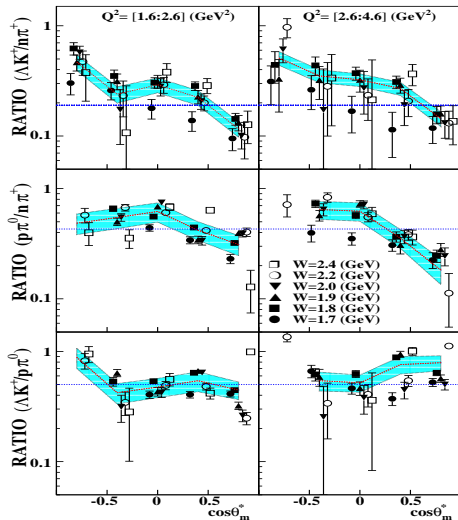
LOI12-15-004, + Additional interesting aspect

- Understanding of physics between $Q^2 = 0$ and $Q^2 > 0 \text{ GeV}^2$
- Already saw in some resonance states in previous presentation
- Another example $K^+\Lambda$ induced polarization [Phys.Rev.C90, 035202 (2014).]



LOI12-15-004, + Additional interesting aspect

- Hadronization in high energy by color flux-tube model
- Strangeness suppression factor shows consistent with high energy results - [Phys.Rev.Lett.**113**, 152004 (2014).]
- Q^2 independence is universal down to pthn point ?



Summary

- We have obtained the differential cross-sections/asymmetries using an exclusive single pion electroproduction data for very wide range of kinematics, **near threshold** $< W < \mathbf{DIS}$ regime, $Q^2 = 1.6-4.5 \text{ GeV}^2$.
- Precision of single pion data from CLAS allows to extract the helicity amplitudes for various resonance states
 $N(1440)1/2^+$, $N(1520)3/2^-$, $N(1535)1/2^-$, $N(1675)5/2^-$,
 $N(1680)5/2^+$, and $N(1710)1/2^+$
- Combined analysis with available and future data on all exclusive meson electroproduction channels at $W > 1.2 \text{ GeV}$ at $Q^2 > 2 \text{ GeV}^2$ within the framework of coupled channel approaches will improve considerably our knowledge on N^* -state electro-couplings.